

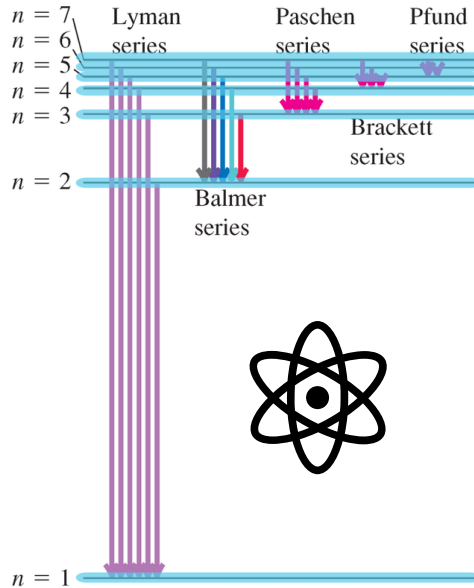
# Unveiling new topological phases with optical fiber networks

Albert Adiyatullin, Corentin Lechevalier, Rabih El Sokhen,  
Clement Evain, François Copie, Stéphane Randoux,  
Pierre Suret, Alberto Amo

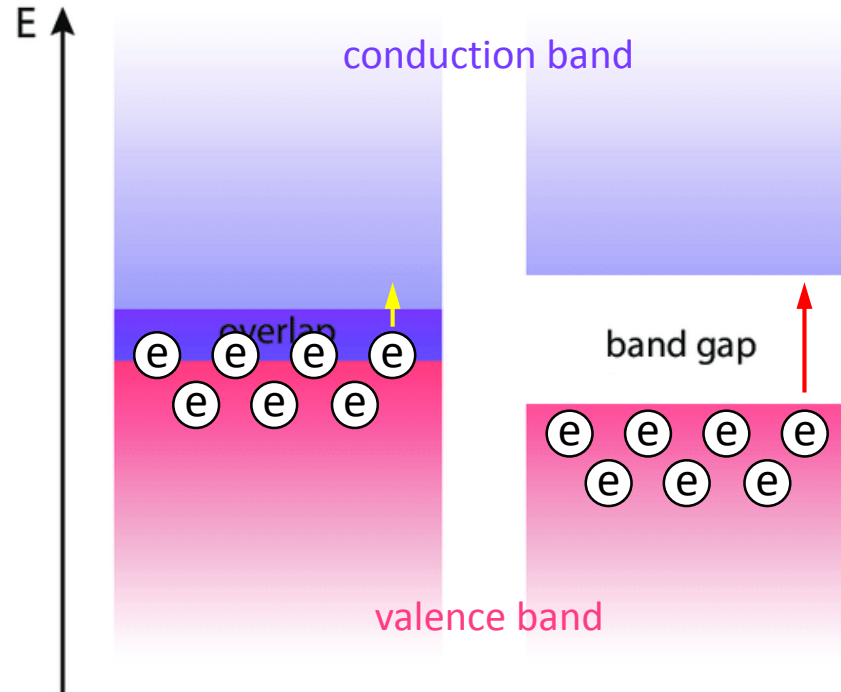
In collaboration with:  
Lavi Upreti (Univ. Würzburg), Pierre Delplace (Univ. Lyon)

- I. Introduction: topology, Floquet driving
- II. Programmable 2D photonic network
- III. Floquet winding metals

## Atom: energy levels



## Crystal: energy bands



CONDUCTORS

INSULATORS

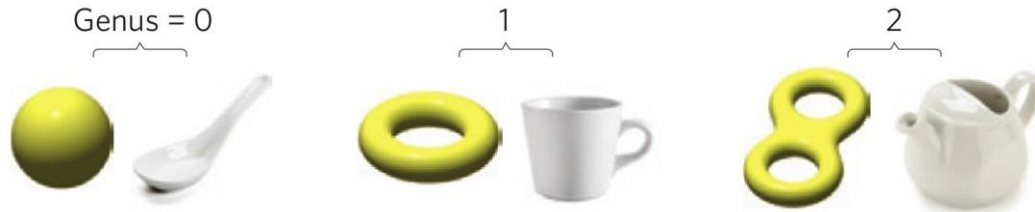
incandescent bulb



LED lamp



## Mathematical example: number of holes



Lu et al., Nature Phot. 2014



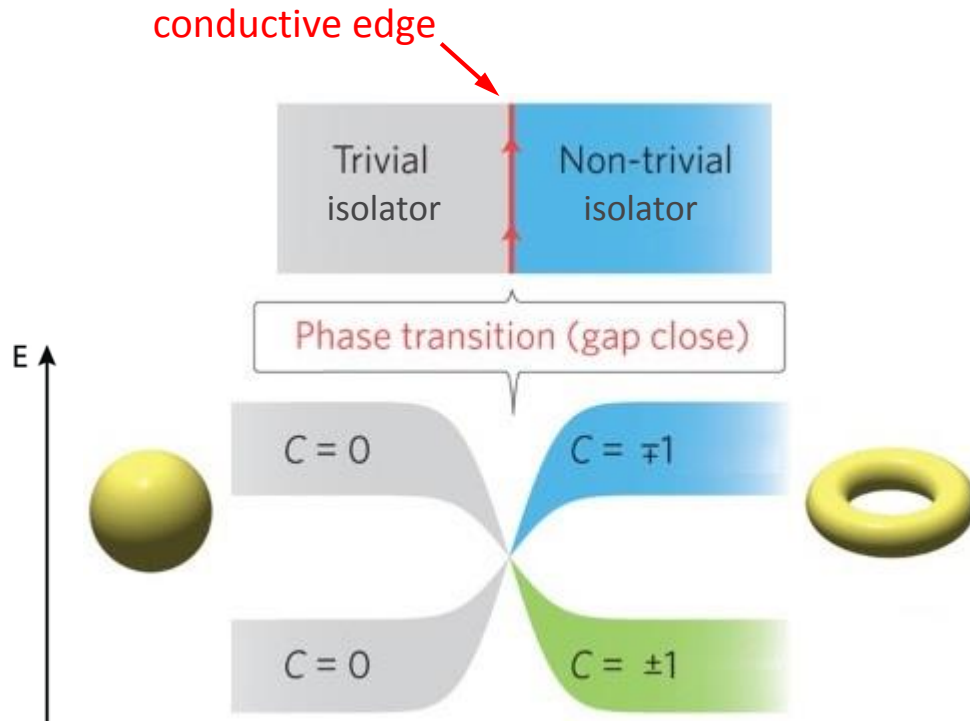
band gap



INSULATORS

- integer number
- can be defined for each band
- immune to perturbations

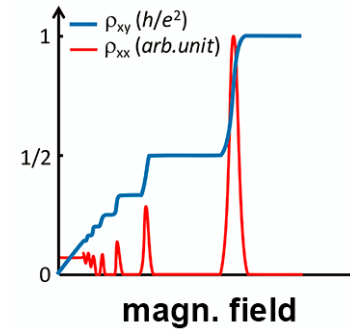
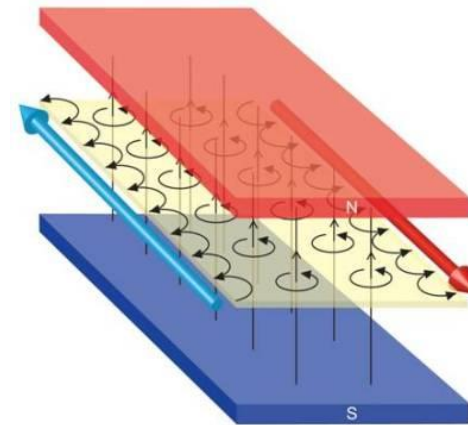
Invariant can change only if the gap closes



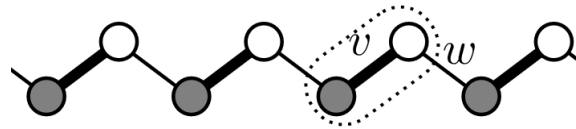
Lu et al., Nature Phot. 2014

topological charge  $\leftrightarrow$  presence of edge states

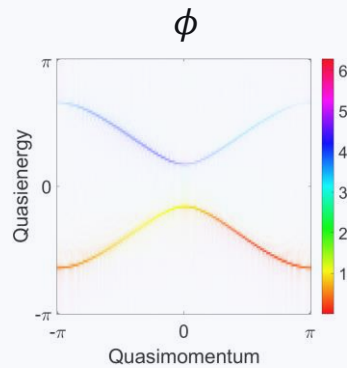
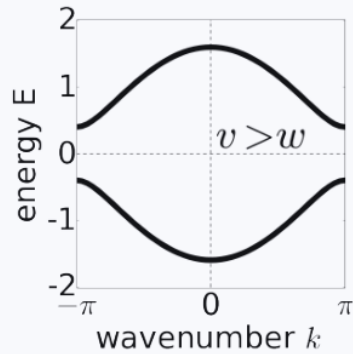
## Quantum Hall effect



Bipartite lattice (2 sites per unit cell):



Trivial:  $v > w$

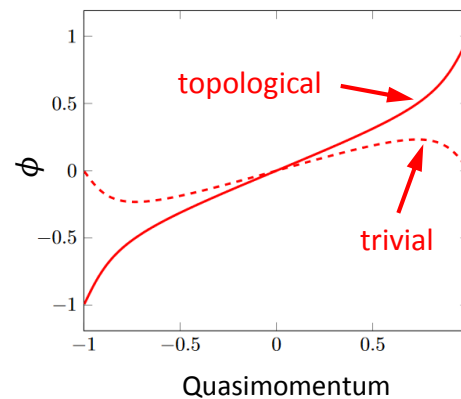


Stationary Schrödinger equation  $\hat{H}\Psi = E\Psi$

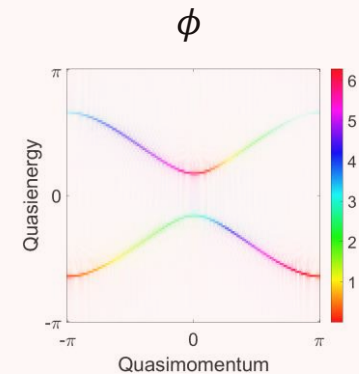
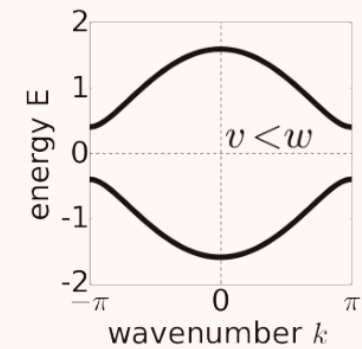
Wavefunctions:  $\Psi_{\pm} = \begin{pmatrix} A \\ B \end{pmatrix}$

Energies:  $E_{\pm}$

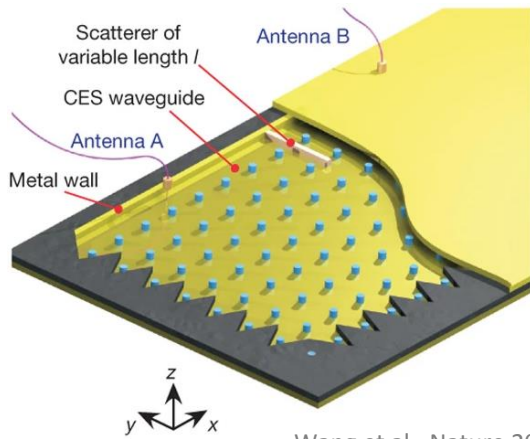
$\phi = \arg(A) - \arg(B)$



Topological:  $v < w$

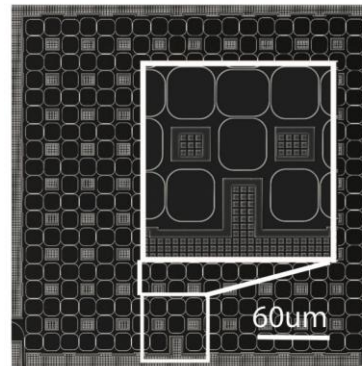


## Microwave resonators



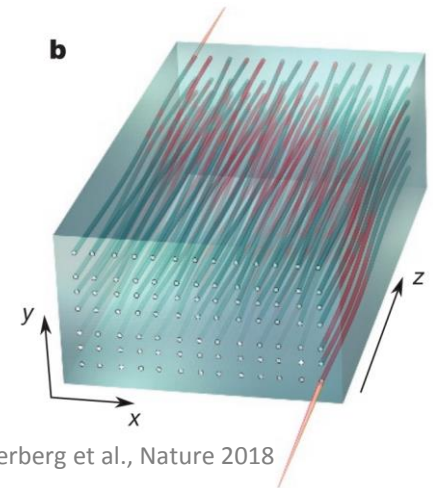
Wang et al., Nature 2009

## Optical resonators



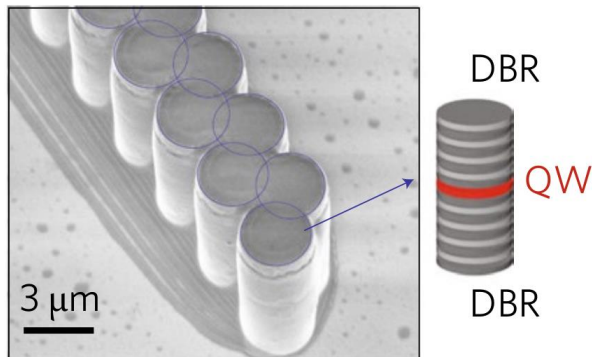
Hafezi et al., Nat. Phot. 2013

## Waveguides



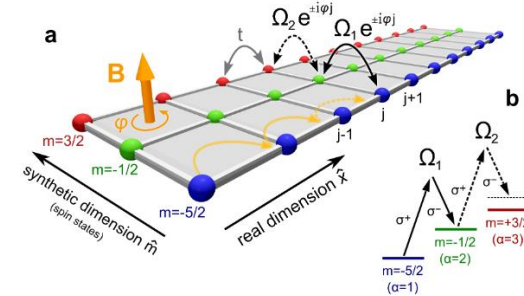
Zilberberg et al., Nature 2018

## Exciton-polaritons



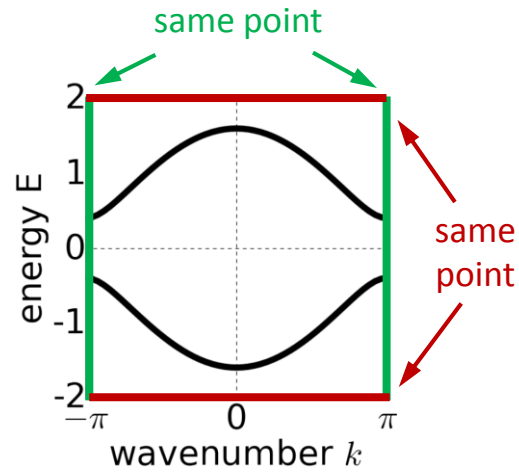
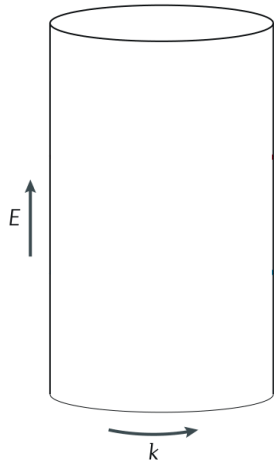
St-Jean et al., Nat. Phot. 2017

## Cold atoms

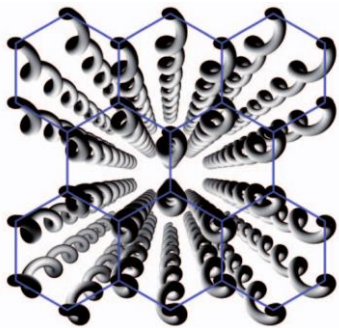
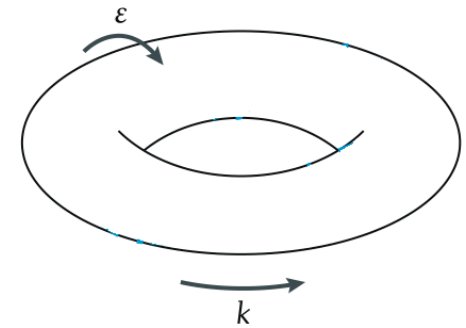


Mancini et al., Science 2015

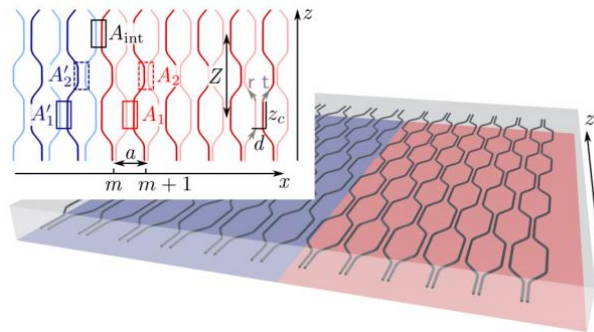
Crystal  
periodicity in *space*



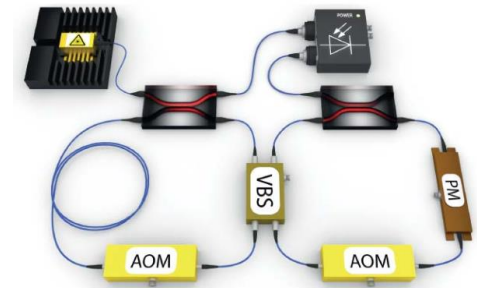
Floquet driving  
periodicity in *time*



Rechtsmann et al., Nature 2013

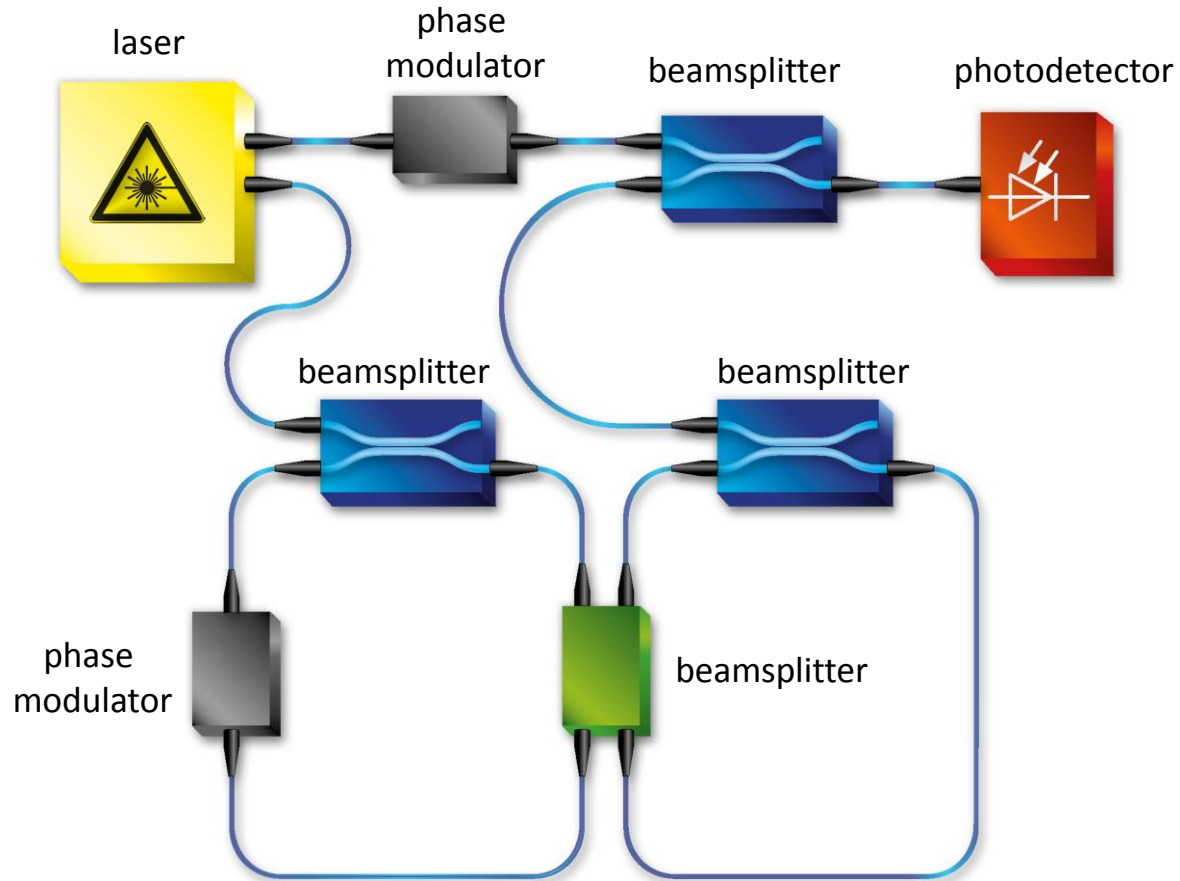


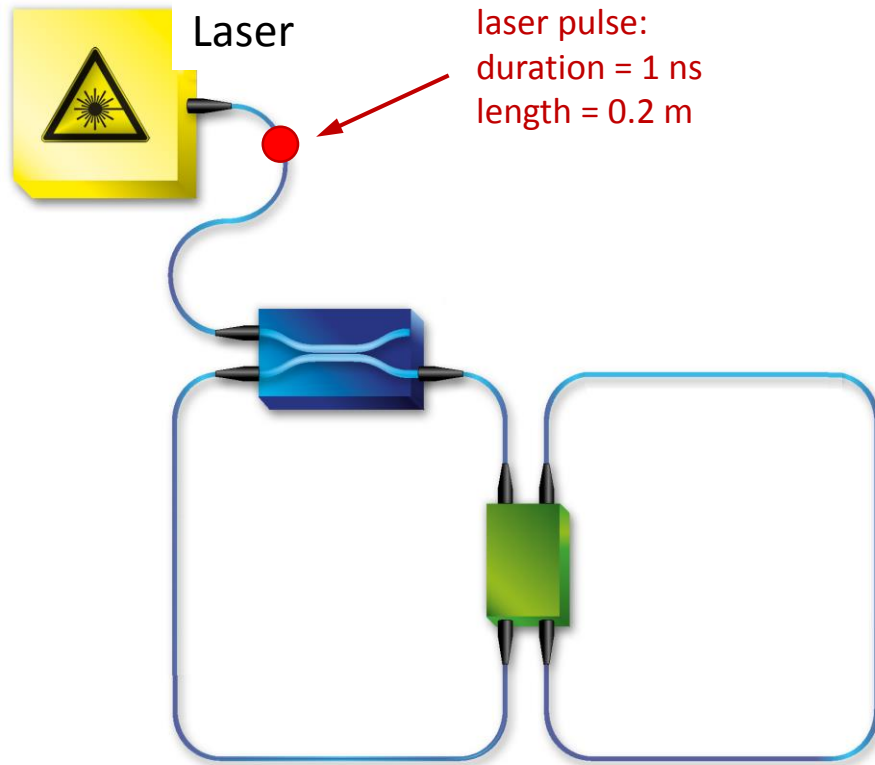
Bellec et al., Europhys. Lett. 2017



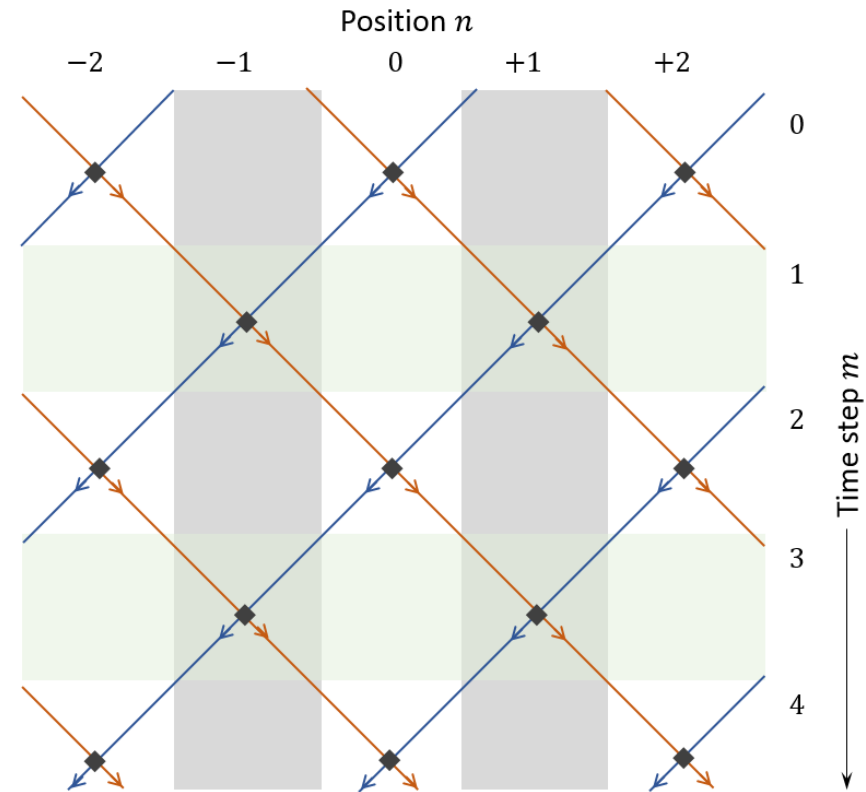
Weidmann et al., Science 2020  
Regensburger et al., Nature 2012

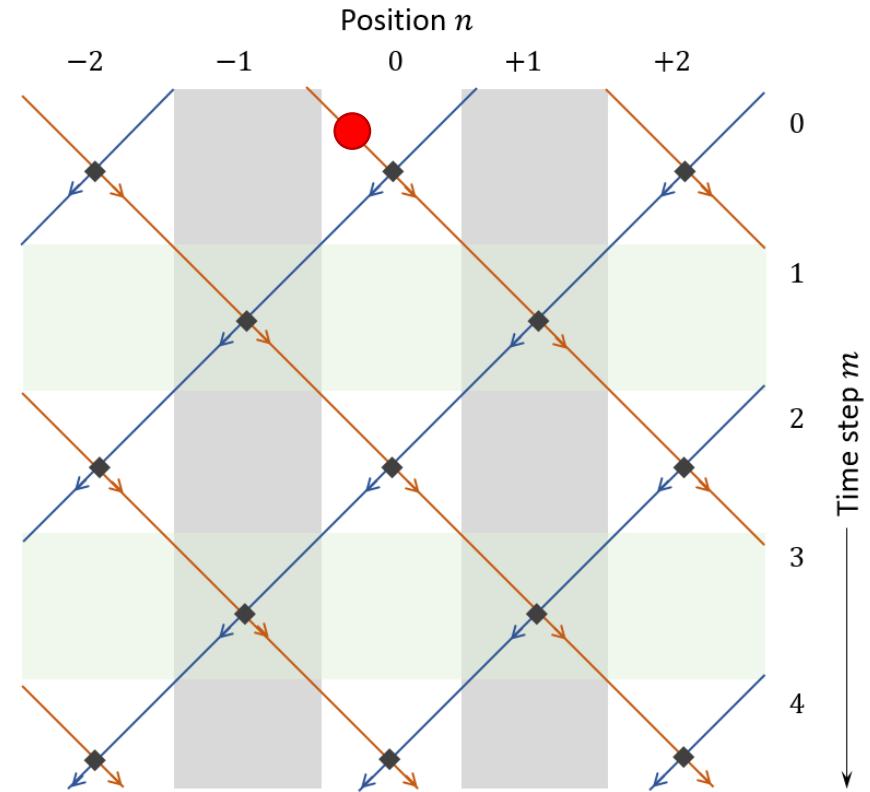
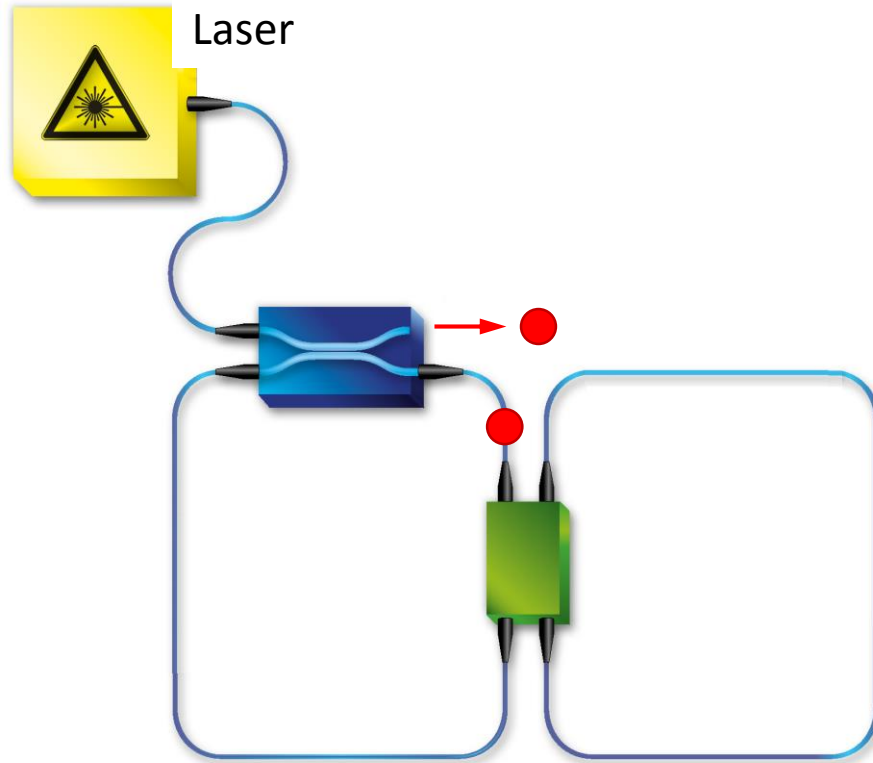




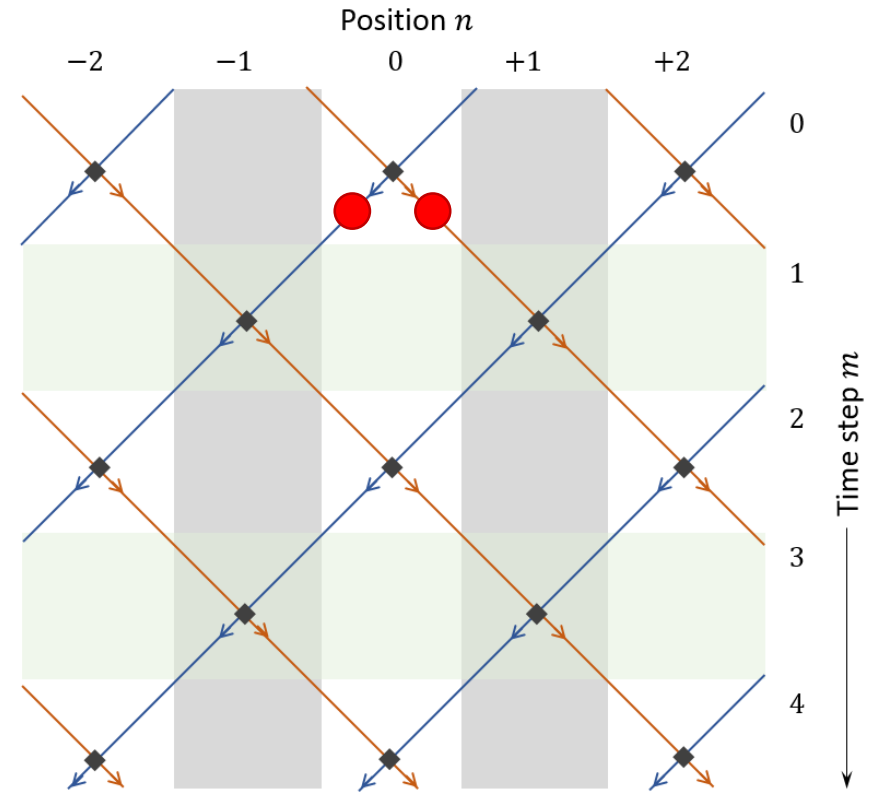
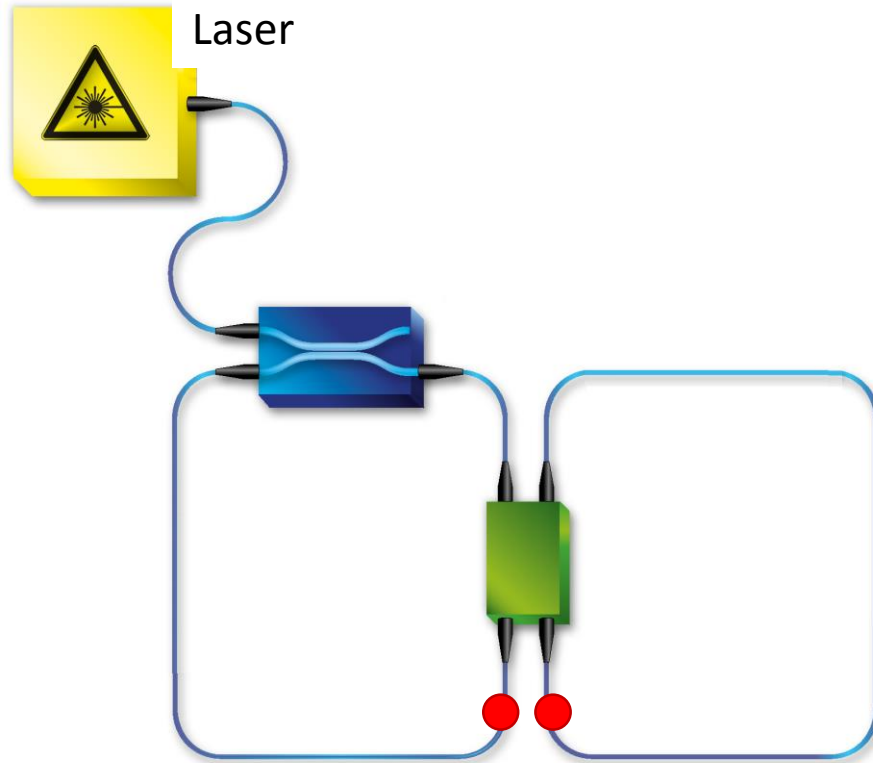


rings  $\approx 40$  m each  
length difference  $\approx 0.5$  m

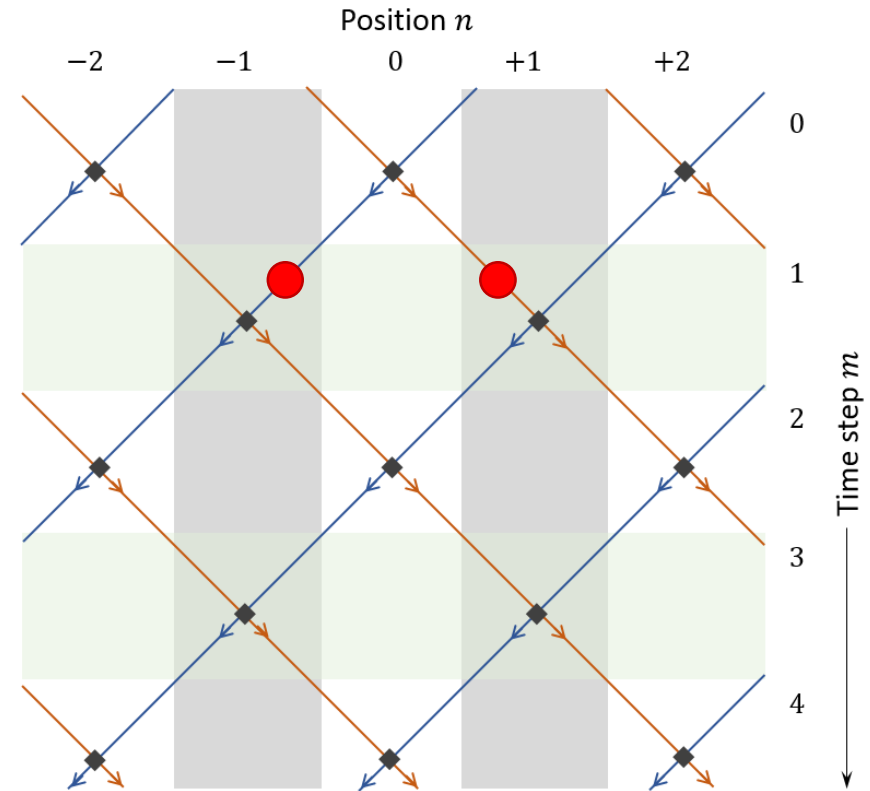
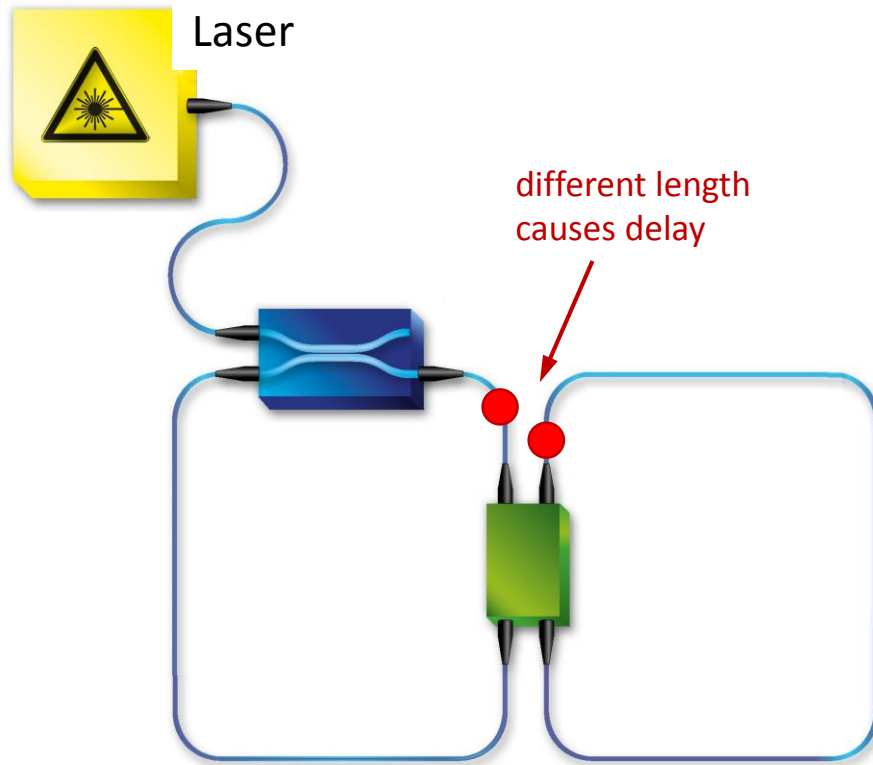




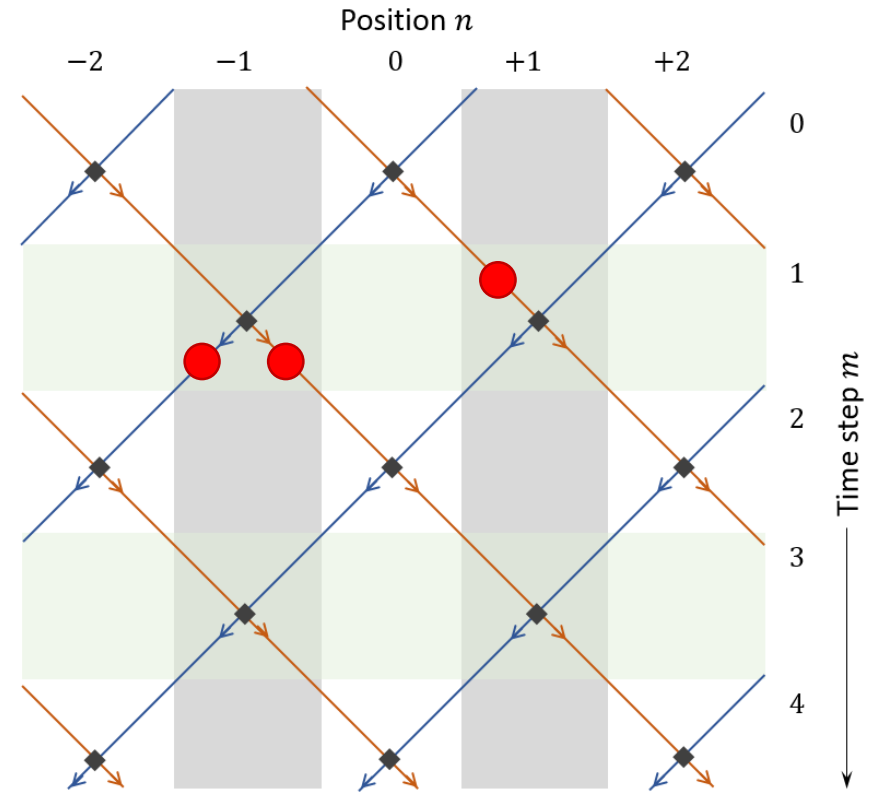
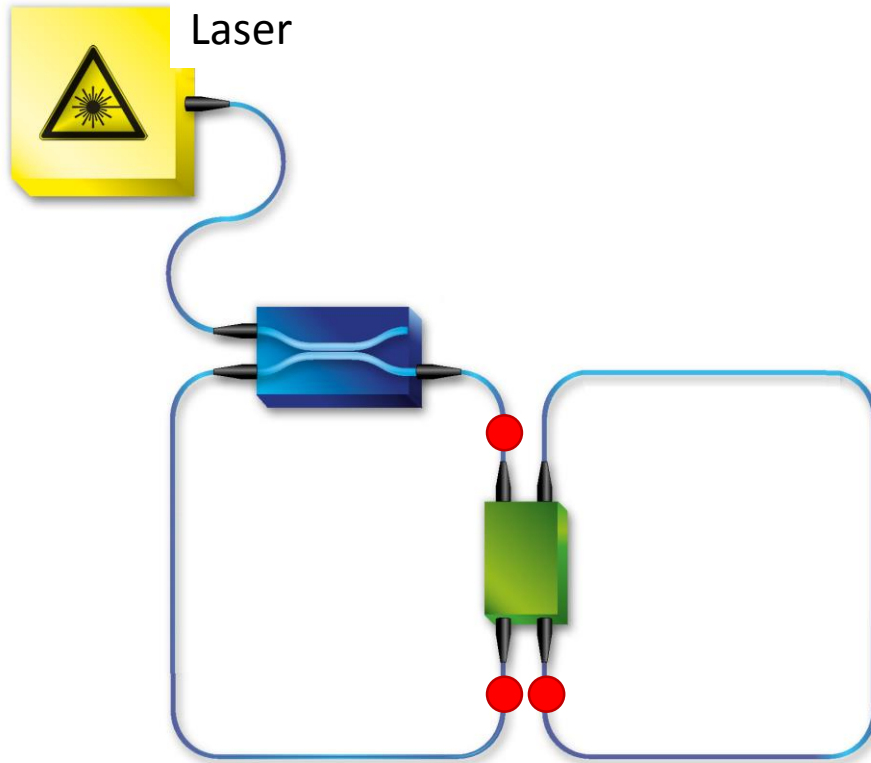
## 1<sup>st</sup> turn



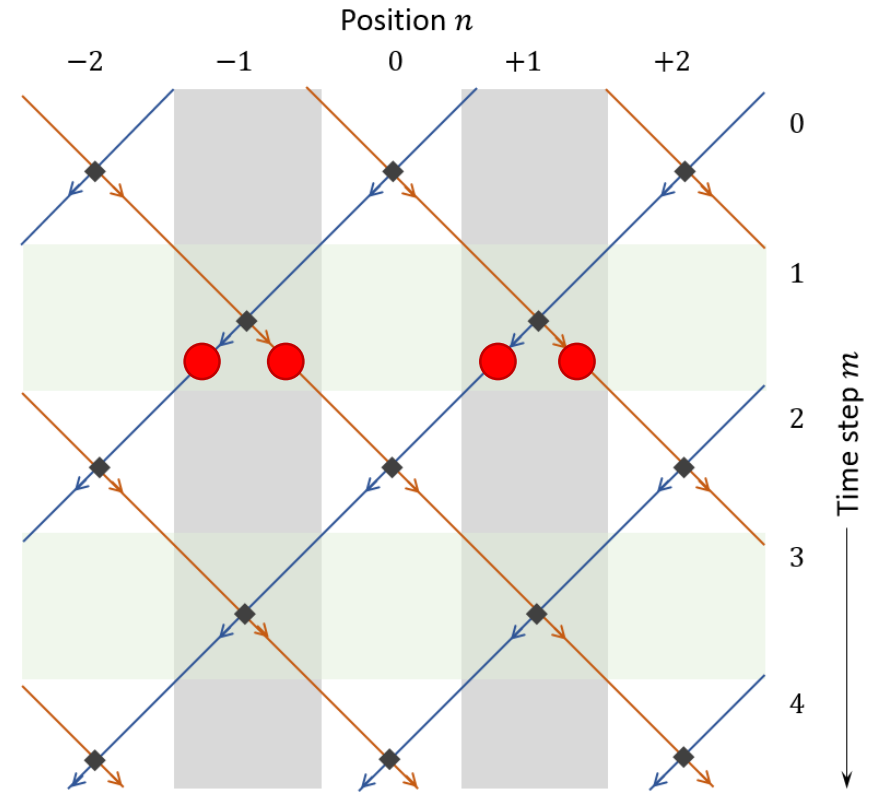
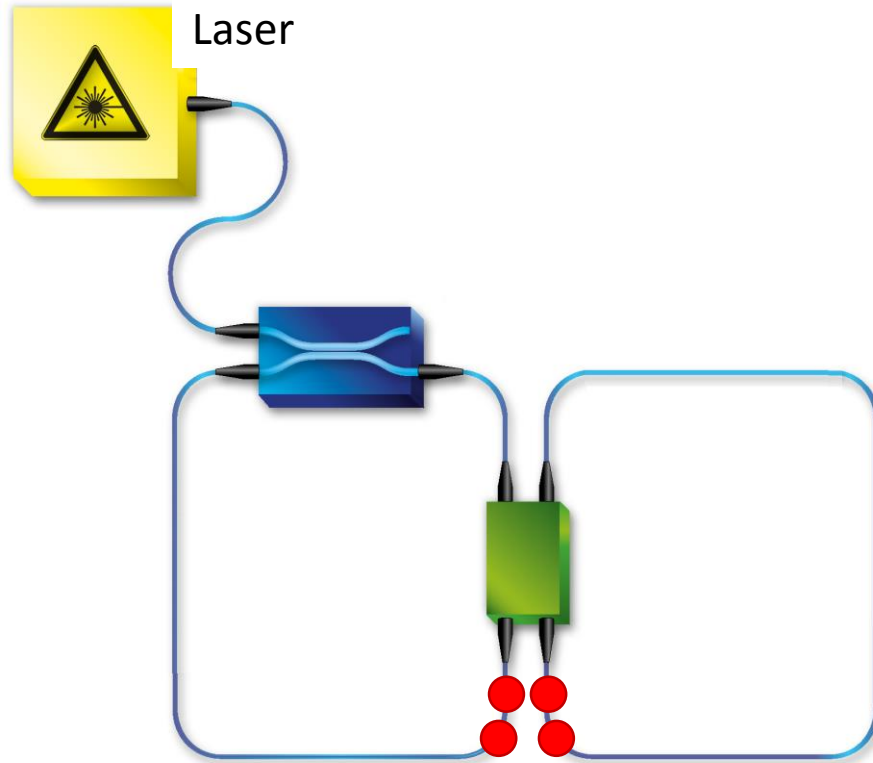
## 1<sup>st</sup> turn



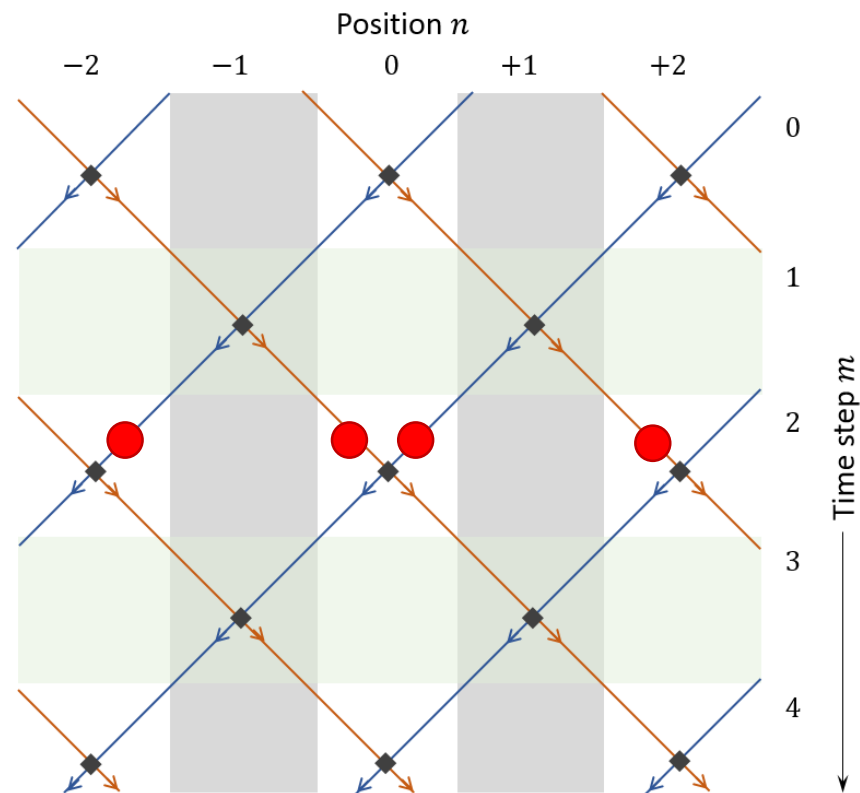
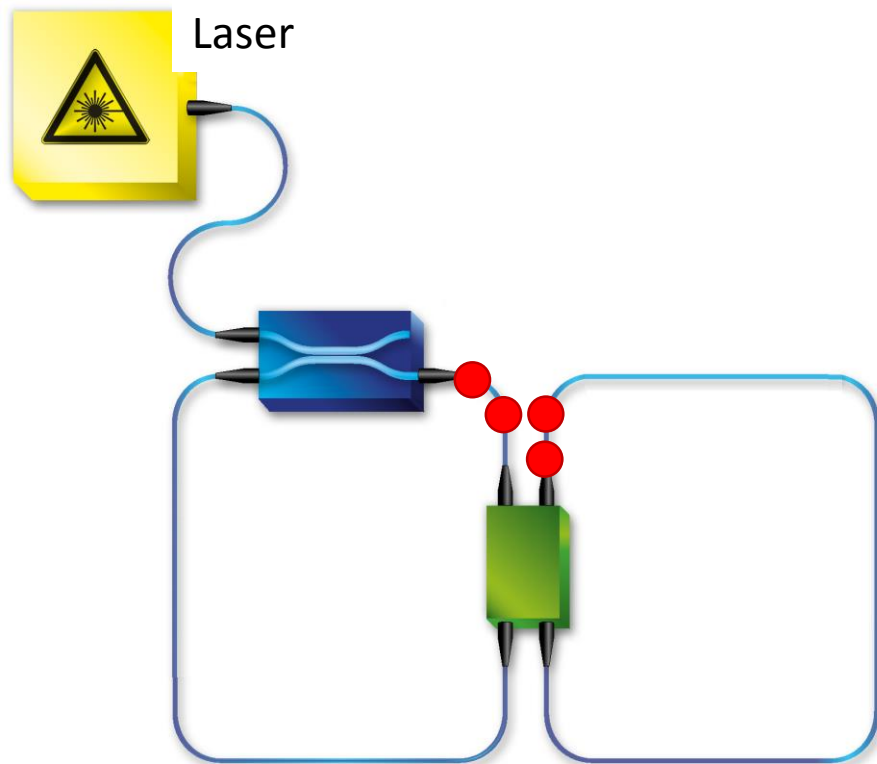
## 2<sup>nd</sup> turn



## 2<sup>nd</sup> turn

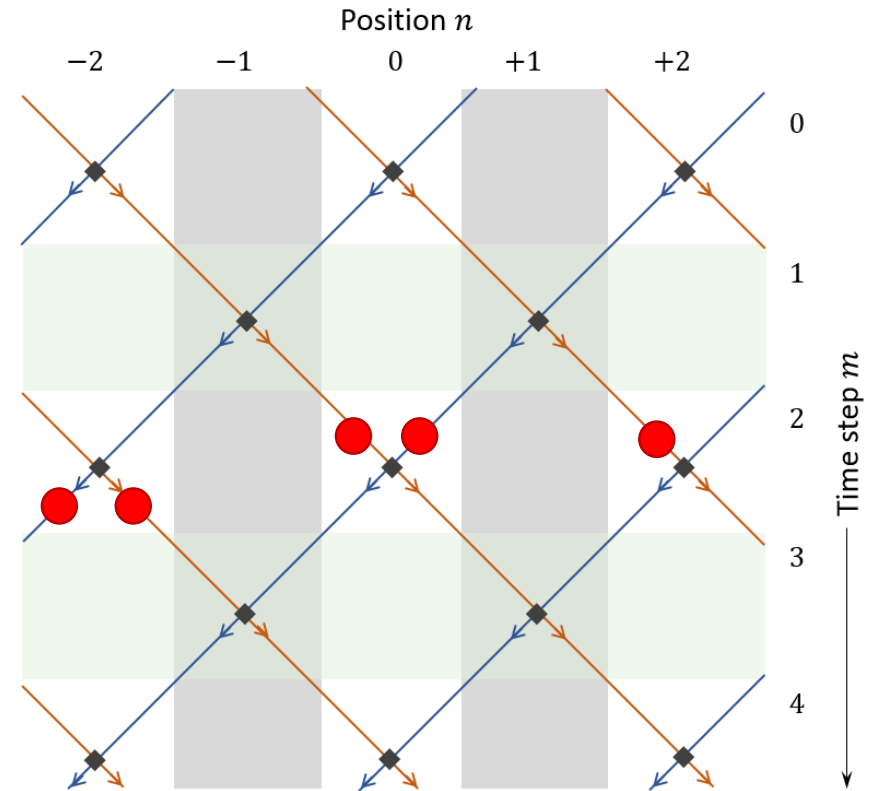
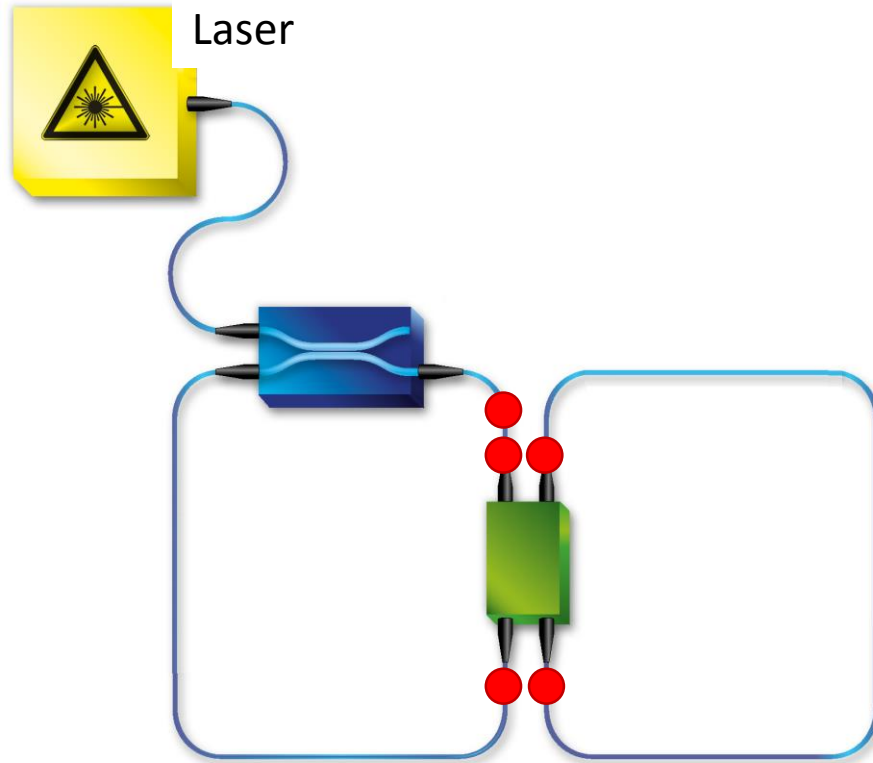


## 2<sup>nd</sup> turn

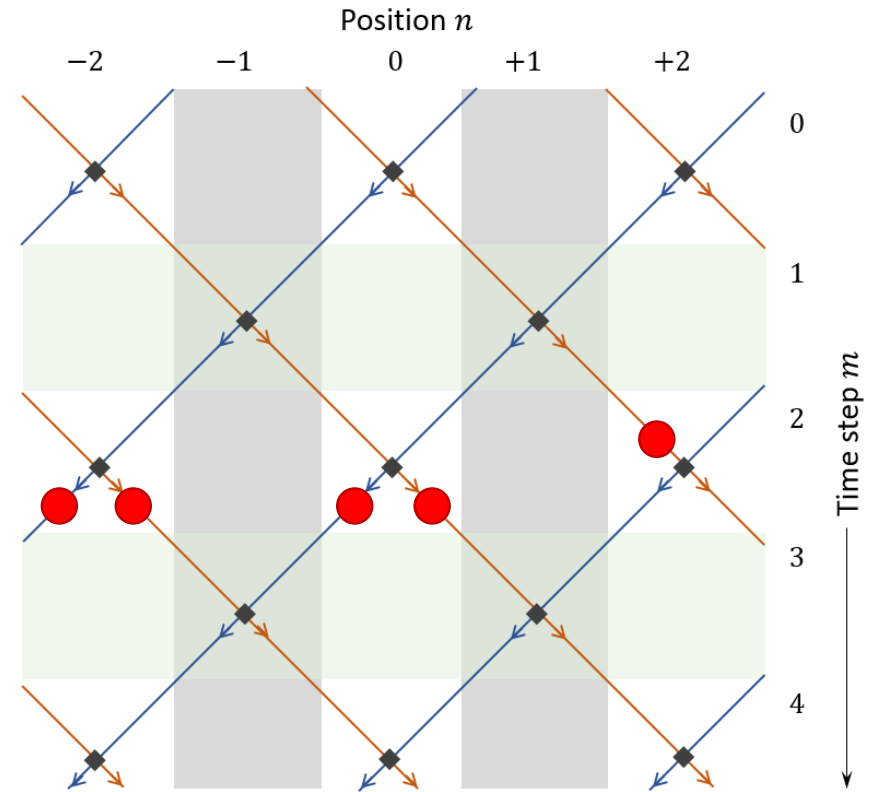
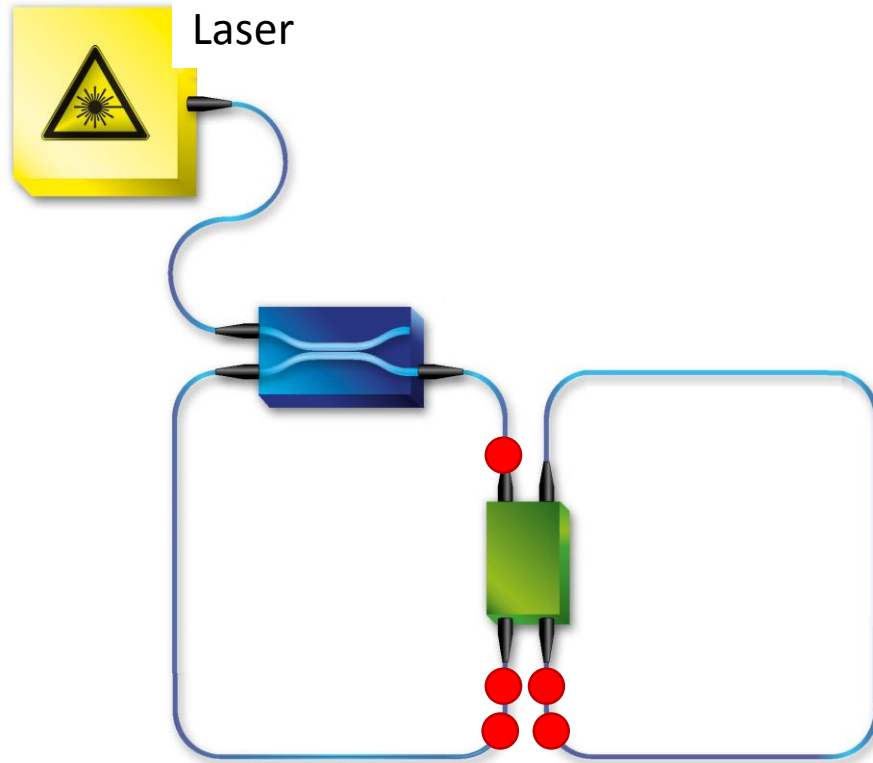




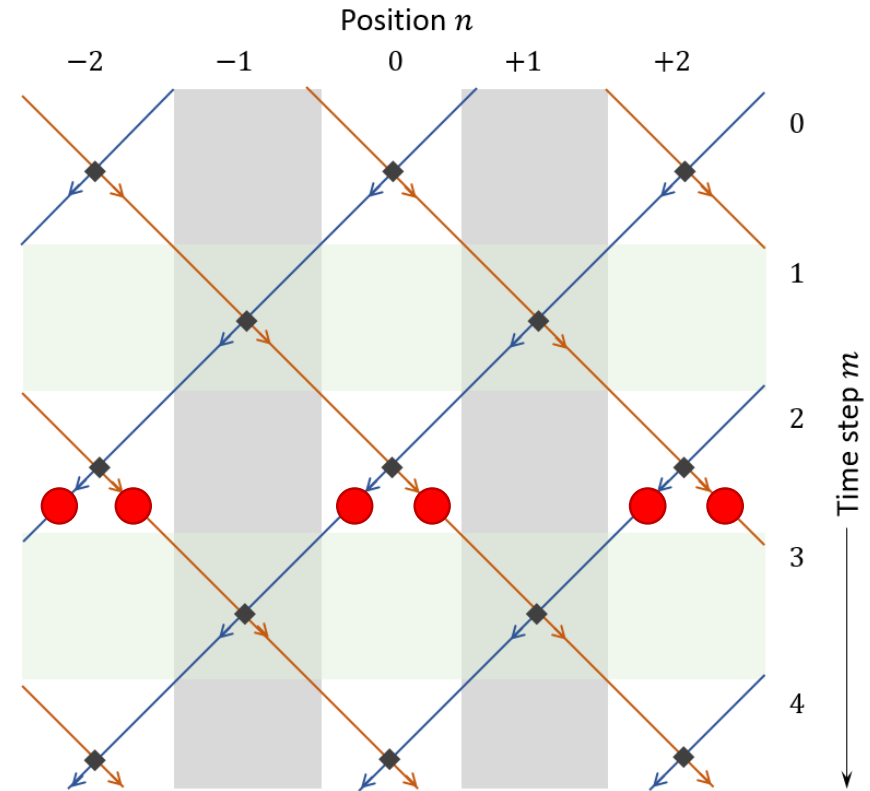
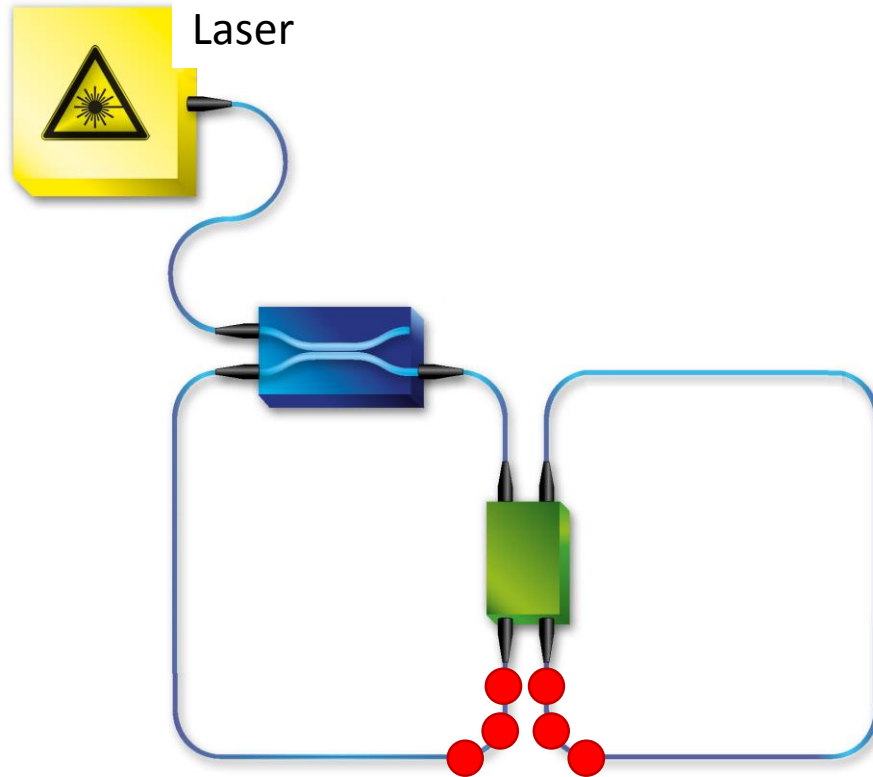
## 3<sup>rd</sup> turn

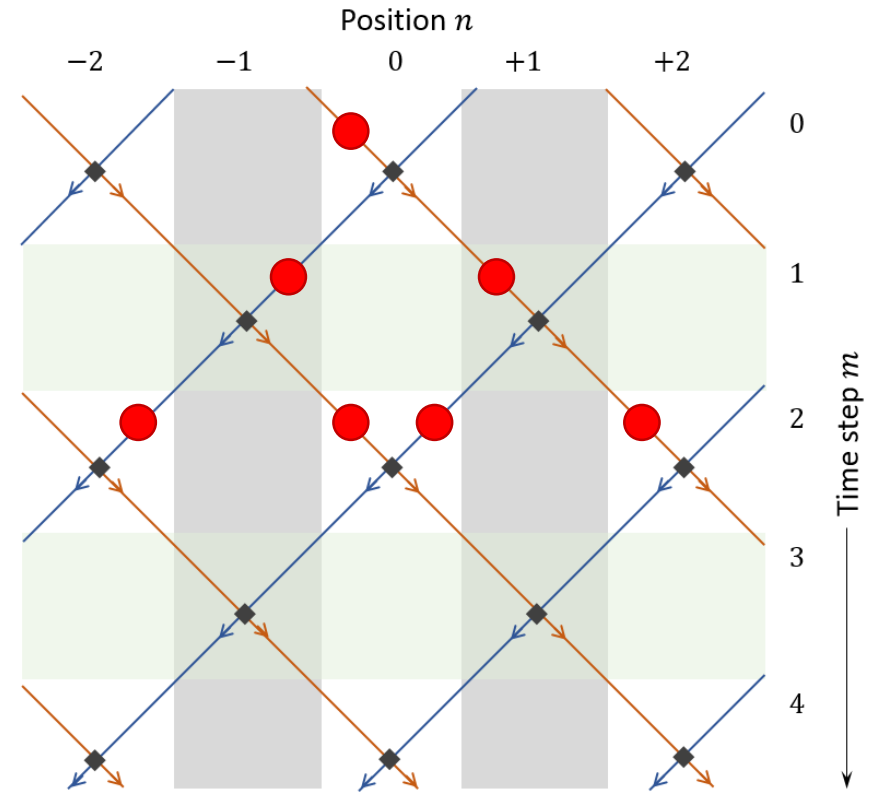
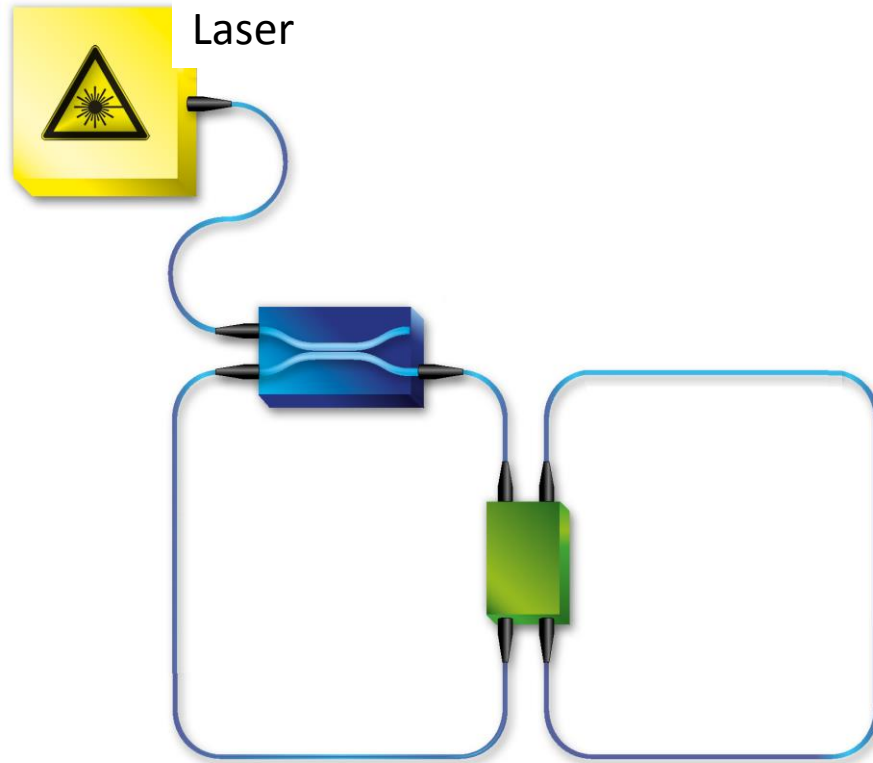


## 3<sup>rd</sup> turn



## 3<sup>rd</sup> turn





$$\alpha_n^{m+1} = \frac{1}{\sqrt{2}} \alpha_{n-1}^m + \frac{1}{\sqrt{2}} i \beta_{n-1}^m$$

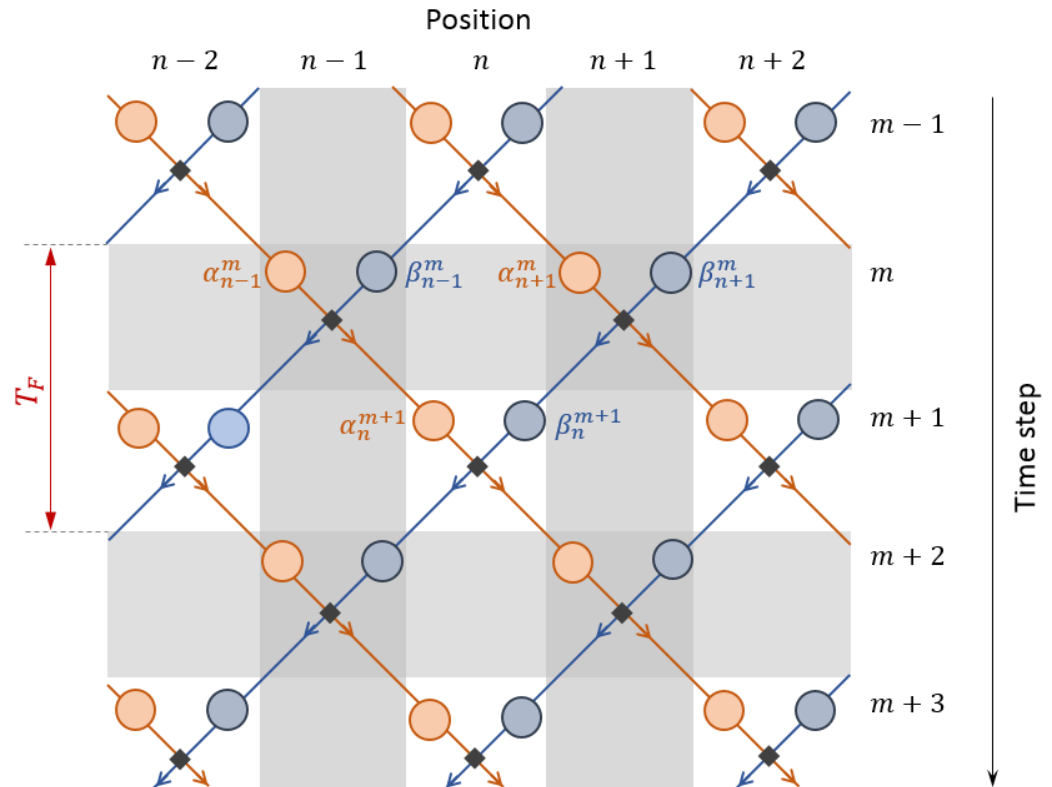
$$\beta_n^{m+1} = \frac{1}{\sqrt{2}} i \alpha_{n+1}^m + \frac{1}{\sqrt{2}} \beta_{n+1}^m$$

Floquet-Bloch ansatz:

$$\begin{pmatrix} \alpha_n^m \\ \beta_n^m \end{pmatrix} = \begin{pmatrix} A \\ B \end{pmatrix} e^{i\frac{kn}{2}} e^{i\frac{Em}{2}}$$

Bands:

$$E = \pm \frac{1}{2} \cos^{-1}(1 - \cos k)$$



$$\alpha_n^{m+1} = \frac{1}{\sqrt{2}} \alpha_{n-1}^m + \frac{1}{\sqrt{2}} i \beta_{n-1}^m$$

$$\beta_n^{m+1} = \frac{1}{\sqrt{2}} i \alpha_{n+1}^m + \frac{1}{\sqrt{2}} \beta_{n+1}^m$$

Floquet-Bloch ansatz:

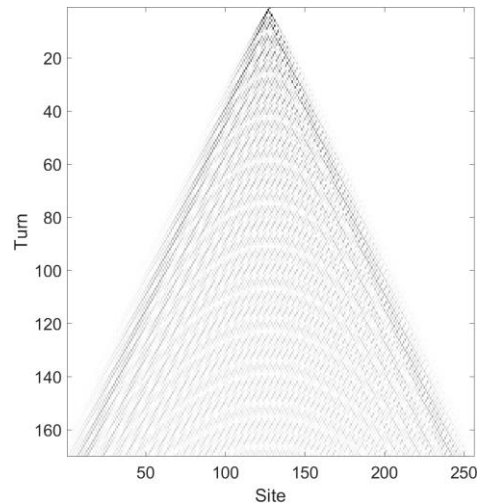
$$\begin{pmatrix} \alpha_n^m \\ \beta_n^m \end{pmatrix} = \begin{pmatrix} A \\ B \end{pmatrix} e^{i \frac{kn}{2}} e^{i \frac{Em}{2}}$$

Bands:

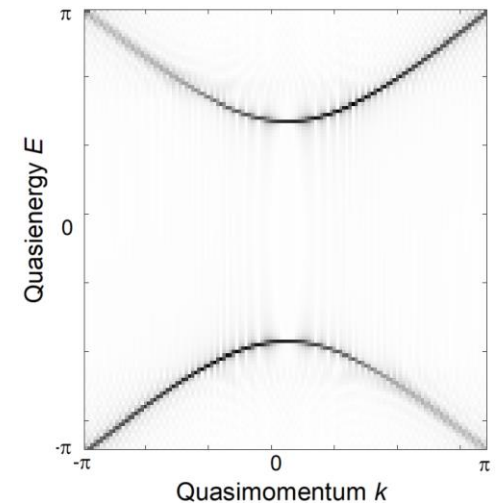
$$E = \pm \frac{1}{2} \cos^{-1}(1 - \cos k)$$

## Simulations

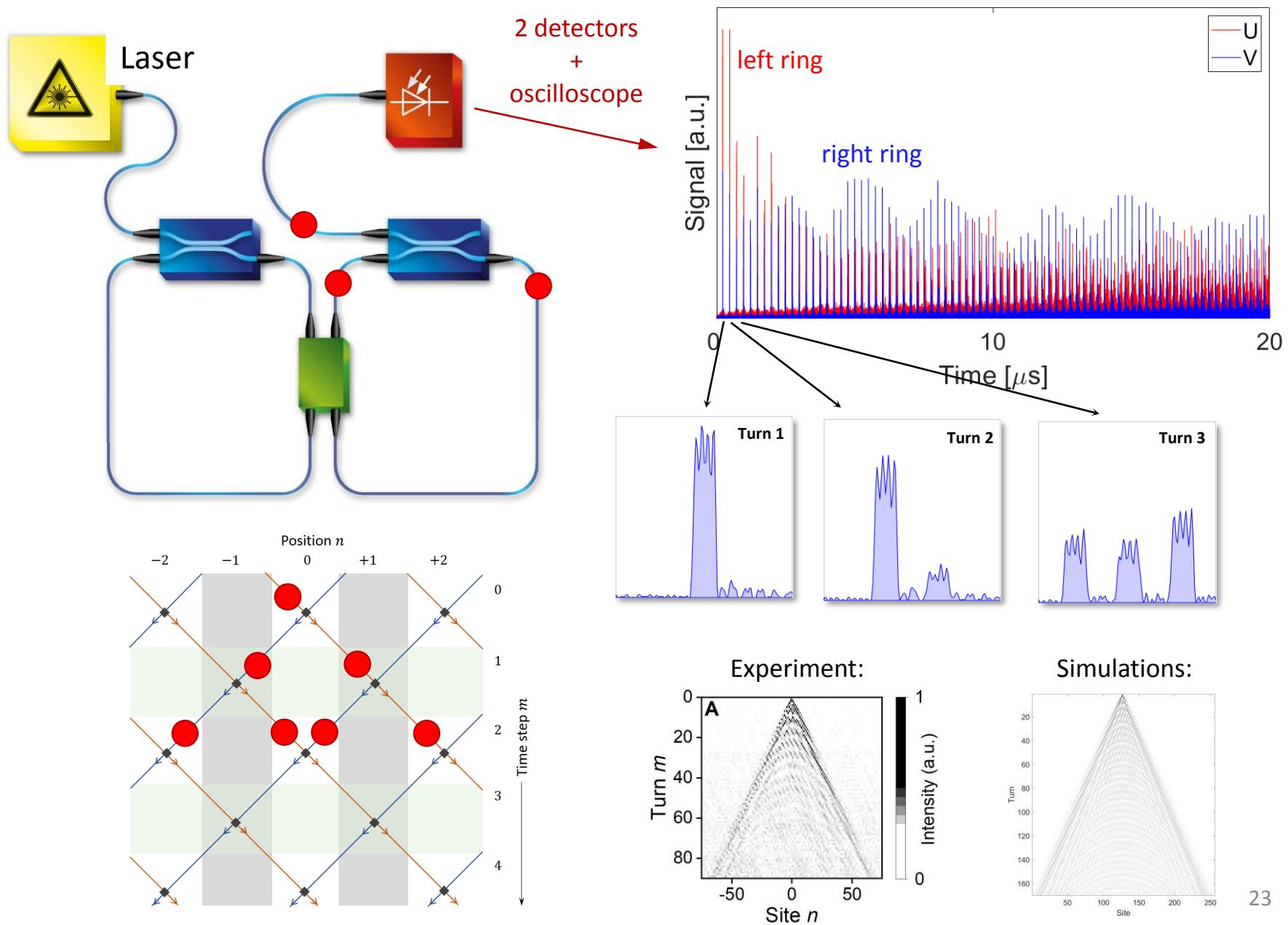
Light walk

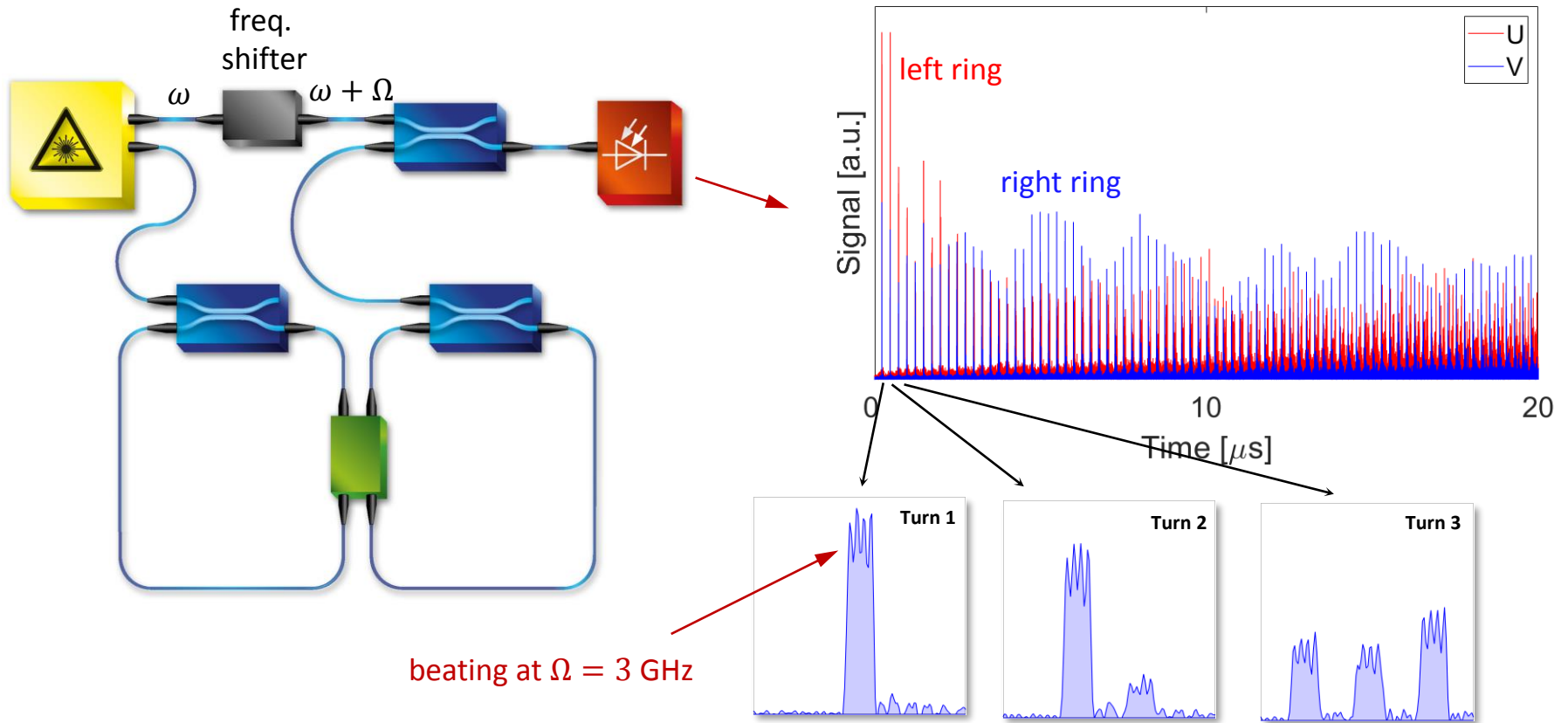


Band structure  
(Fourier transform of the walk)



# Measuring the intensity



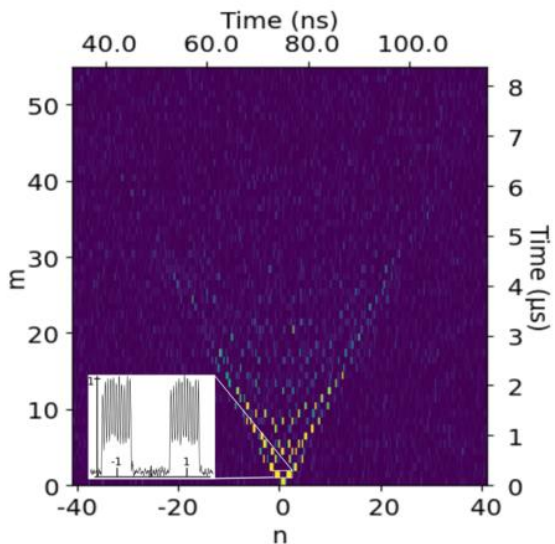


Beating:

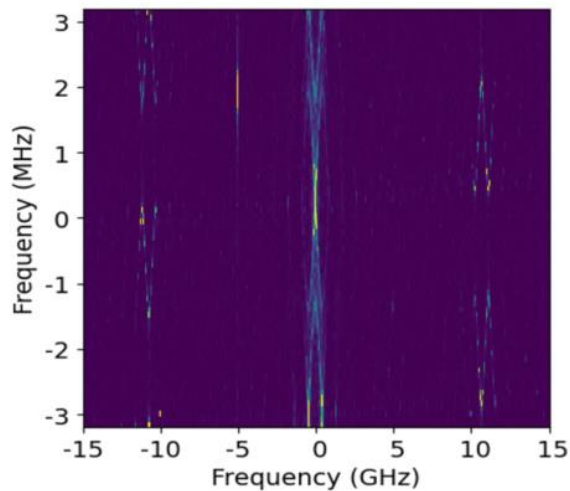
$$I = |A_s e^{i\varphi_s} e^{i\omega t} + A_{LO} e^{i(\omega+\Omega)t}|^2 = A_s^2 + A_{LO}^2 + 2A_s A_{LO} \cos(\Omega t + \varphi_s)$$



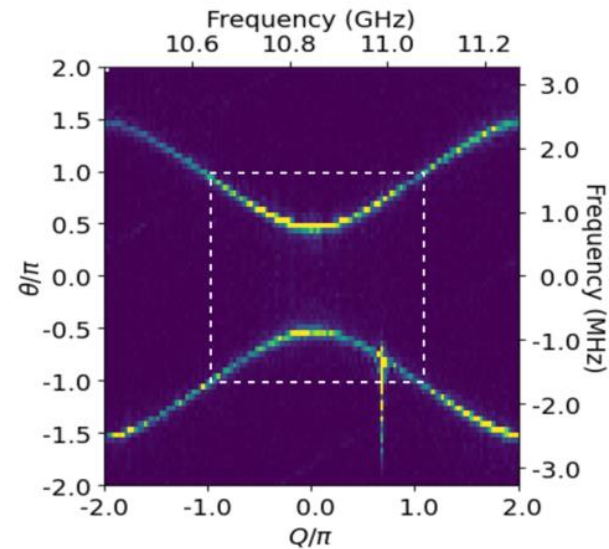
Real space



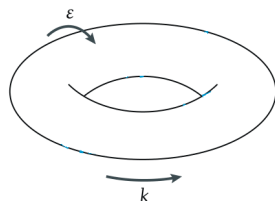
Reciprocal space  
(Fourier transform)



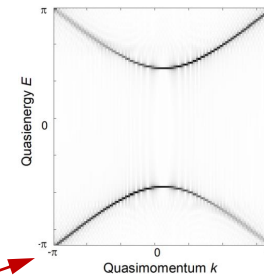
Zoom in at LO frequency



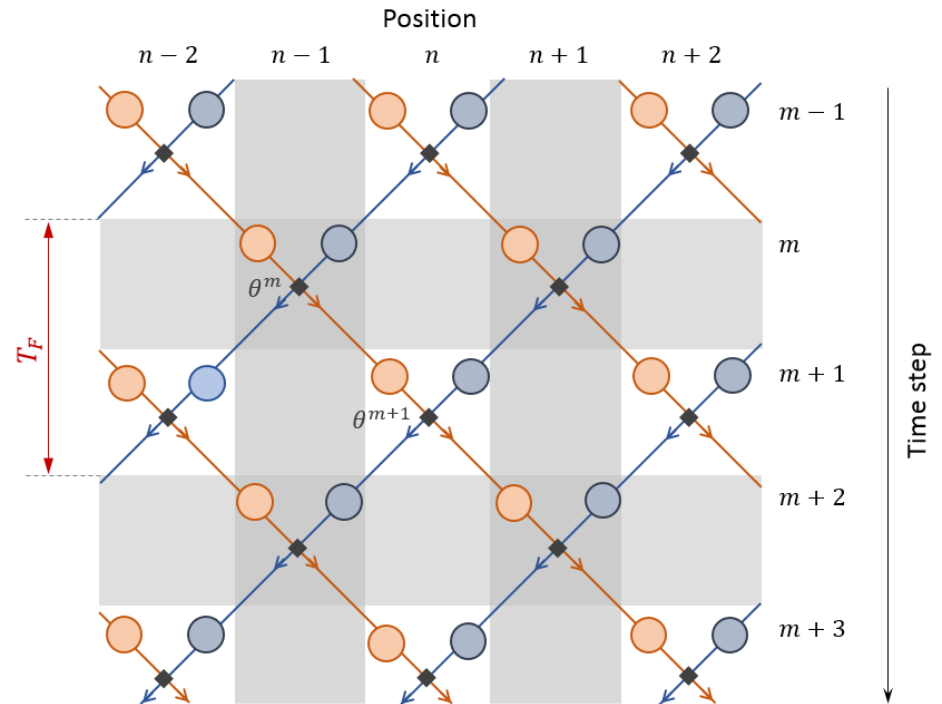
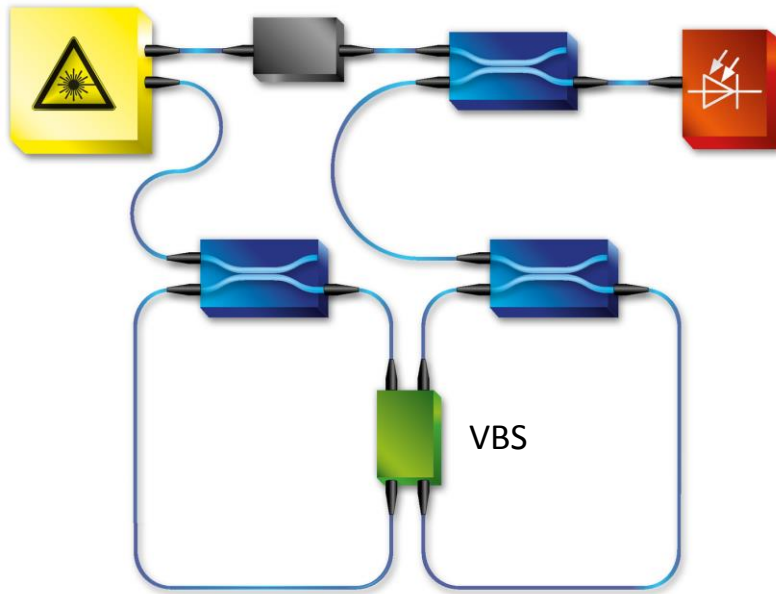
Lechevalier et al., Commun. Phys. 2021



Simulations:



2 touching bands

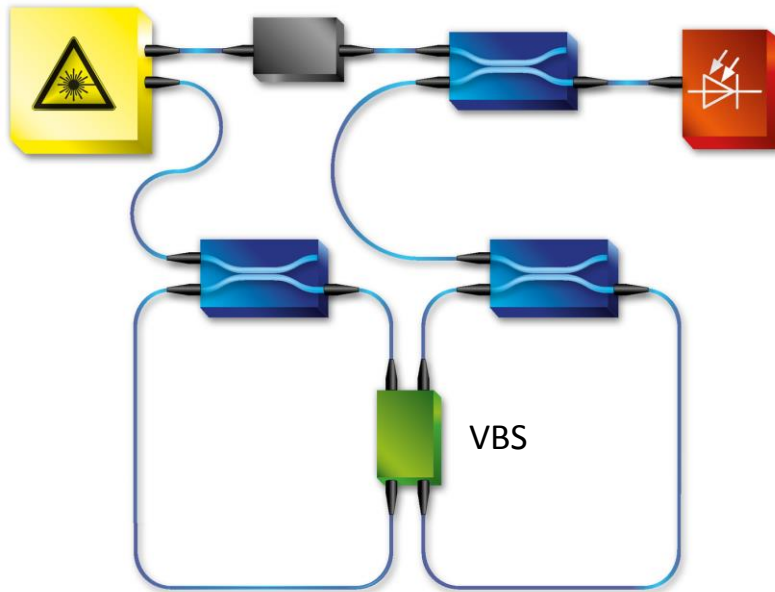


$$\alpha_n^{m+1} = \cos \theta_m \alpha_{n-1}^m + i \sin \theta_m \beta_{n-1}^m$$

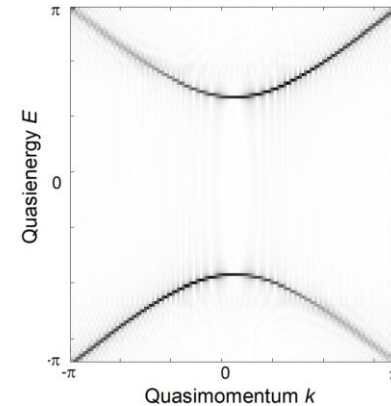
$$\beta_n^{m+1} = i \sin \theta_m \alpha_{n+1}^m + \cos \theta_m \beta_{n+1}^m$$

e.g.  $\theta_m = \pi/4$  means  $\sin \theta_m = \cos \theta_m = \frac{1}{\sqrt{2}}$   
as for 50:50 beamsplitter

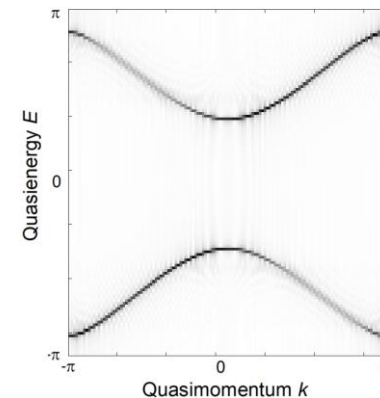
We use 2-step driving:  $\theta_1$  on odd steps,  $\theta_2$  on even steps



$\theta_1 = \theta_2 = \pi/4$  touching bands



$\theta_1 = \pi/4$   
 $\theta_2 = \pi/4 - 0.4$  gap opening

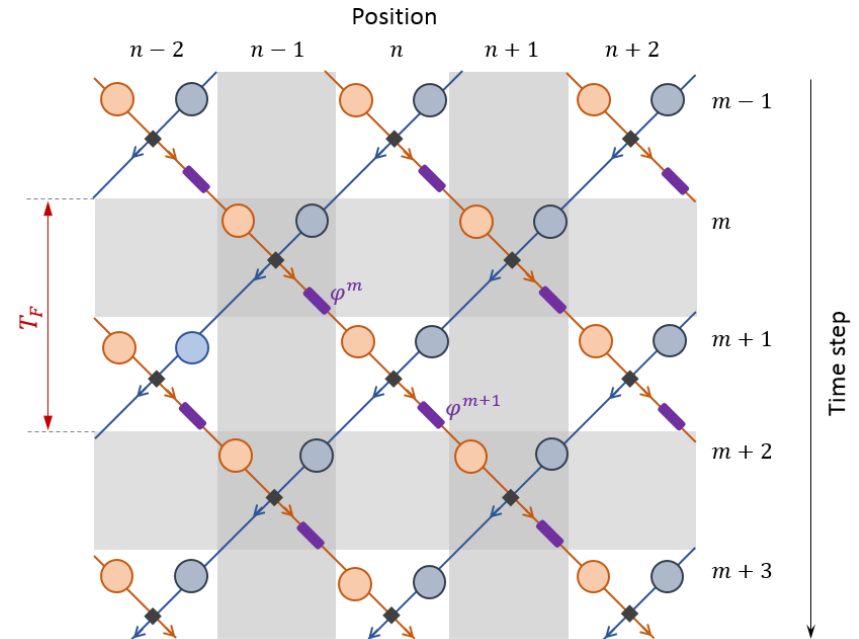
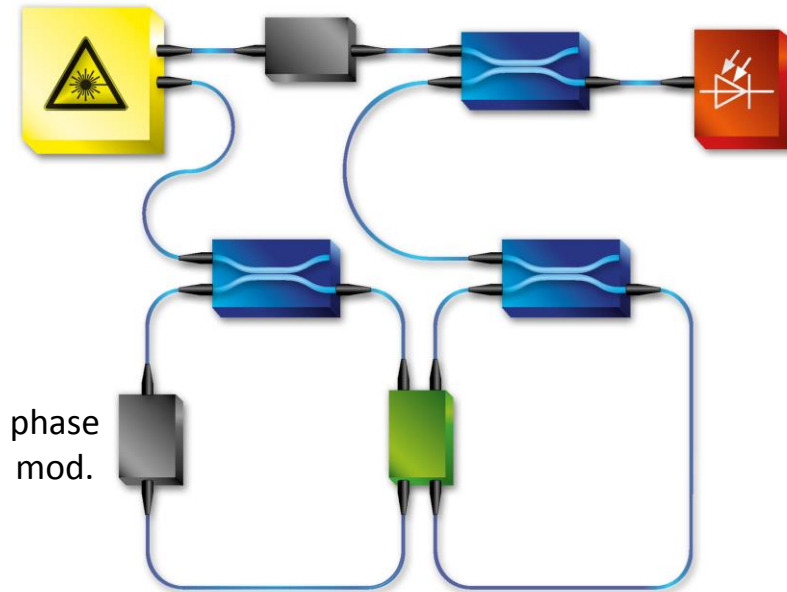


$$\alpha_n^{m+1} = \cos \theta_m \alpha_{n-1}^m + i \sin \theta_m \beta_{n-1}^m$$

$$\beta_n^{m+1} = i \sin \theta_m \alpha_{n+1}^m + \cos \theta_m \beta_{n+1}^m$$

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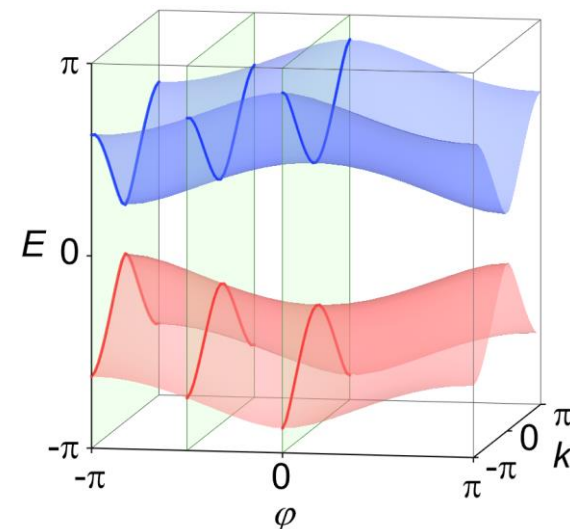
$$\alpha_n^{m+1} = (\cos \theta_m \alpha_{n-1}^m + i \sin \theta_m \beta_{n-1}^m) e^{i\varphi_m}$$

$$\beta_n^{m+1} = i \sin \theta_m \alpha_{n+1}^m + \cos \theta_m \beta_{n+1}^m$$

Again, 2 steps:  $\varphi_1 = c_1 \varphi$  and  $\varphi_2 = c_2 \varphi$

$c_1$  and  $c_2$  are integers

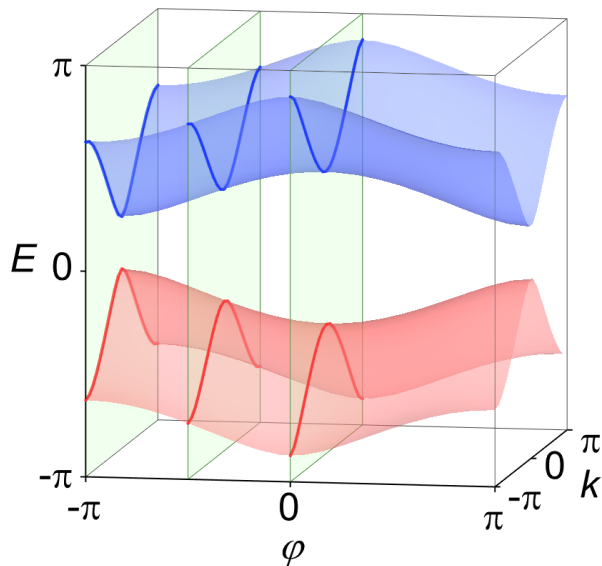
$\varphi \in [-\pi, \pi]$  as a new dimension



Let's introduce integers  $K = (c_1 + c_2)/2$  so that  $\varphi_1 = (K + \Delta)\varphi$   
 $\Delta = (c_1 - c_2)/2$  so that  $\varphi_2 = (K - \Delta)\varphi$

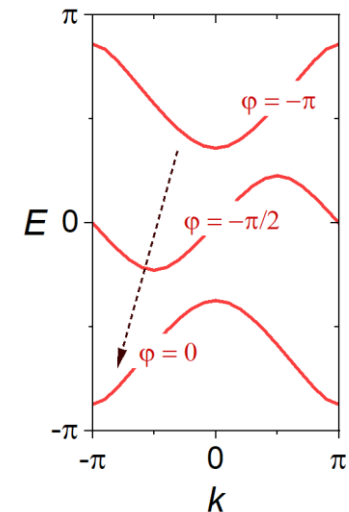
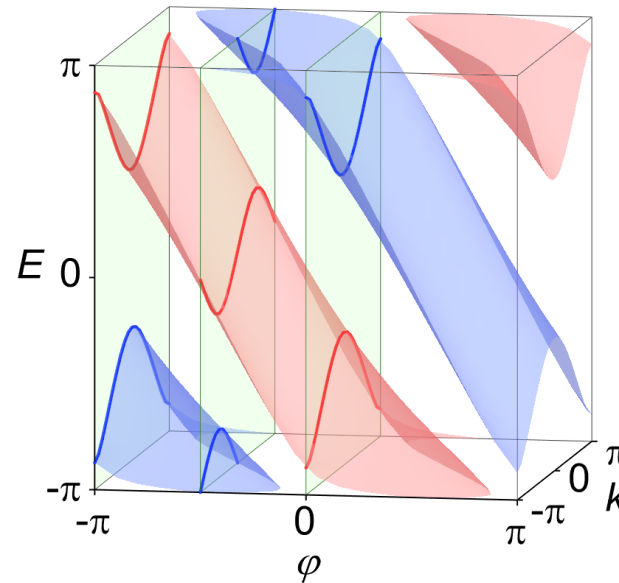
$$K = 0$$

$$\varphi_1 + \varphi_2 = 0$$



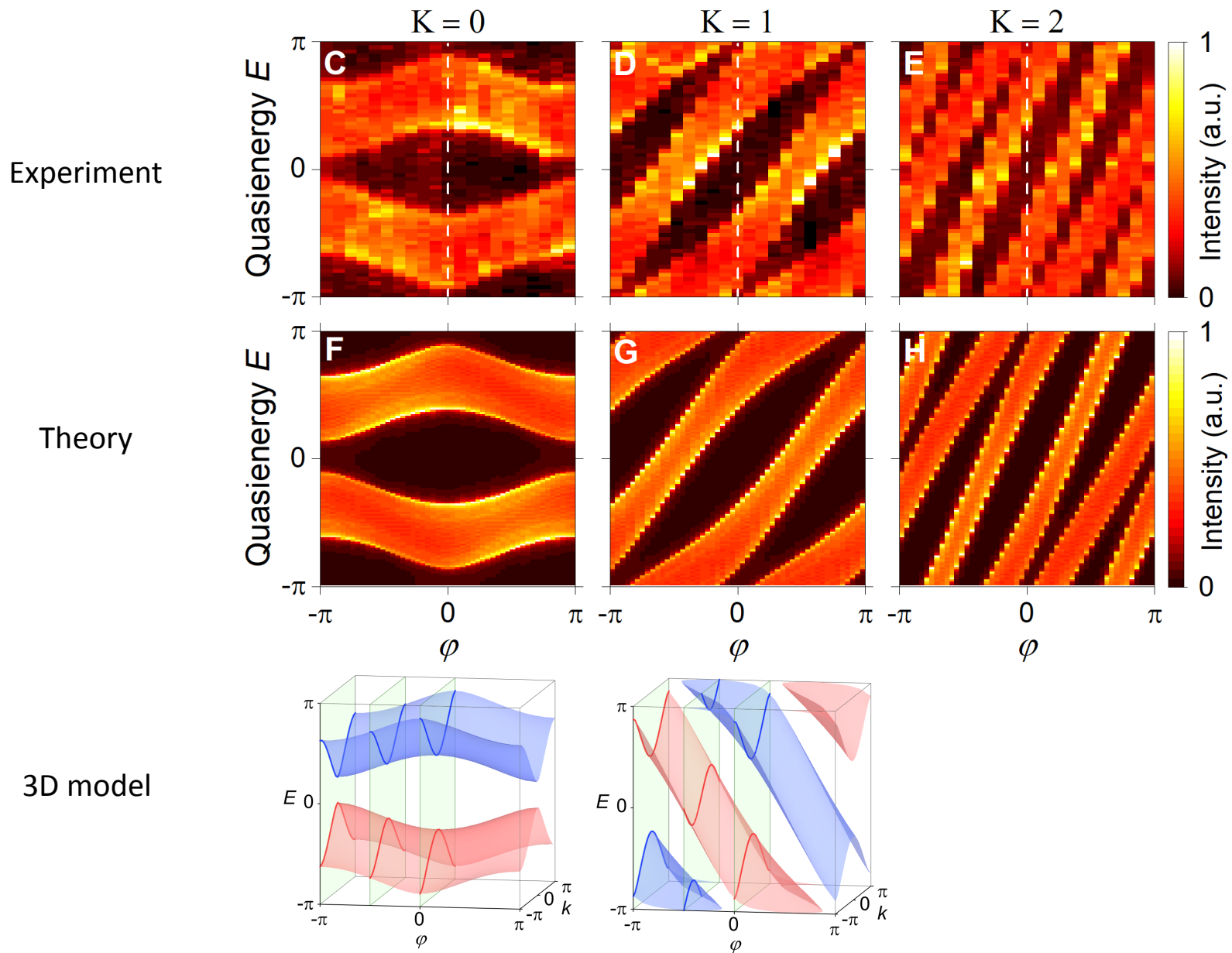
$$K = 1$$

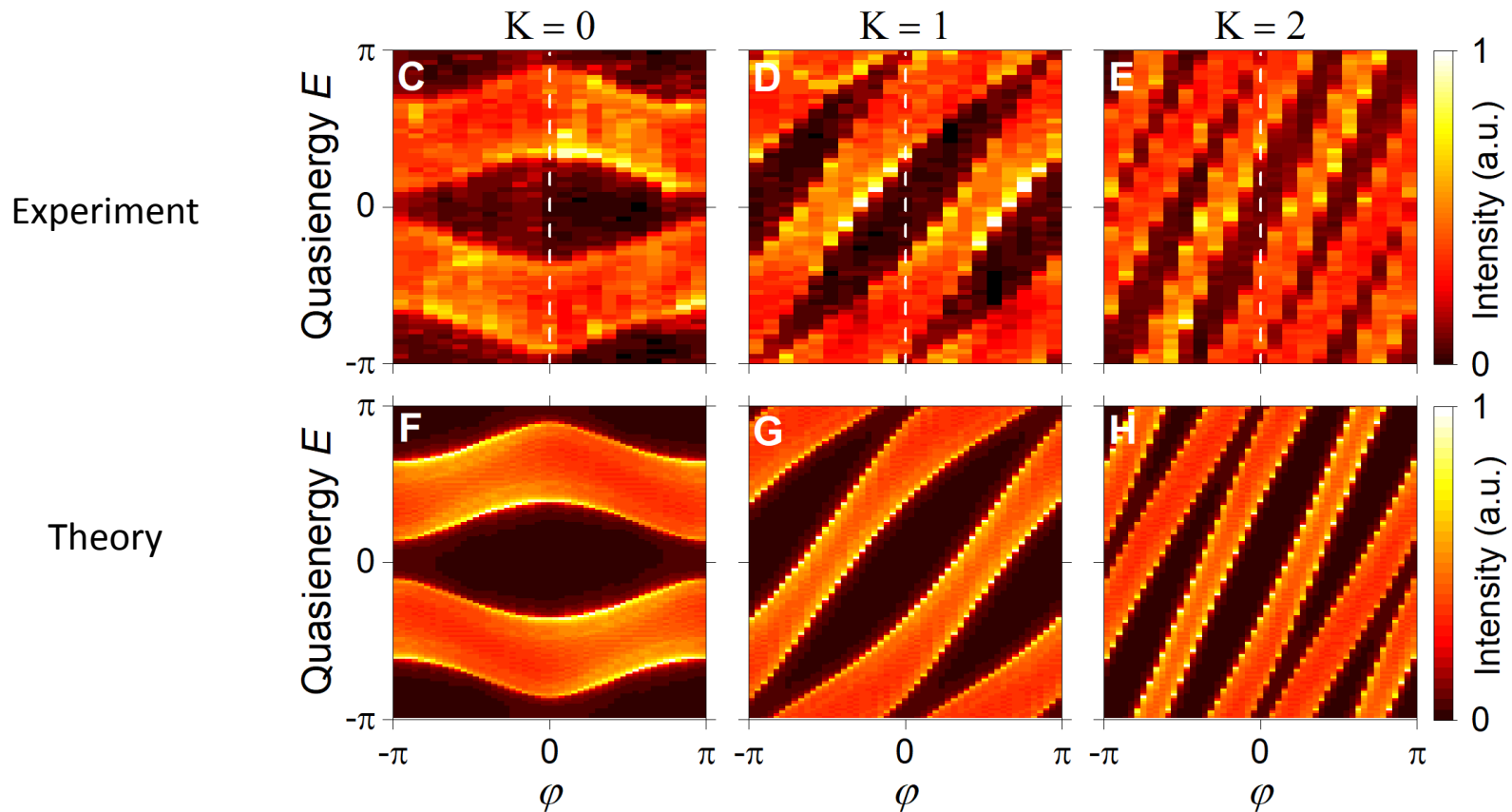
$$\varphi_1 + \varphi_2 \neq 0$$



Floquet winding metal

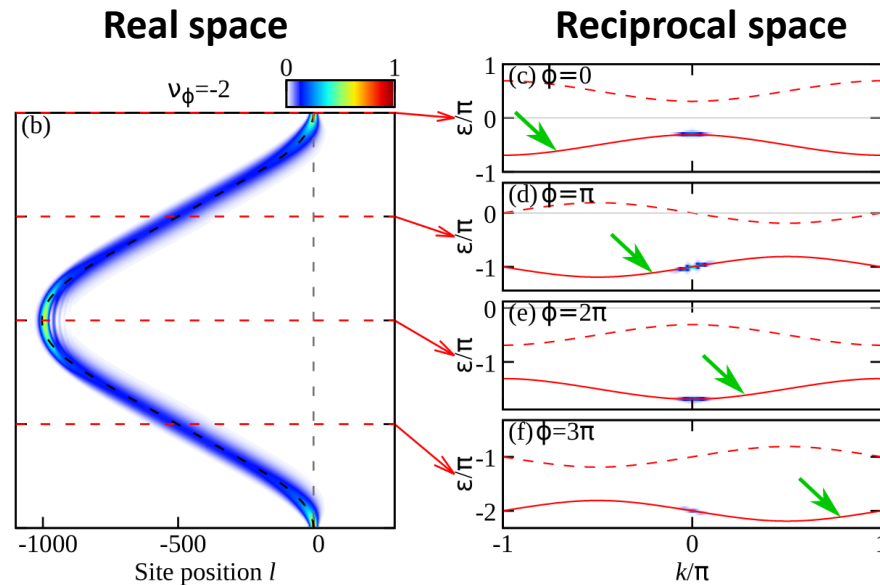
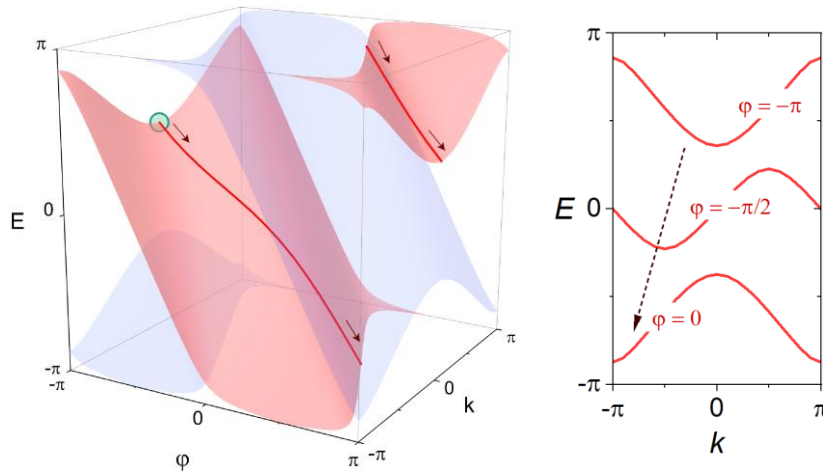
see Upreti et al., Phys. Rev. Lett. 2020





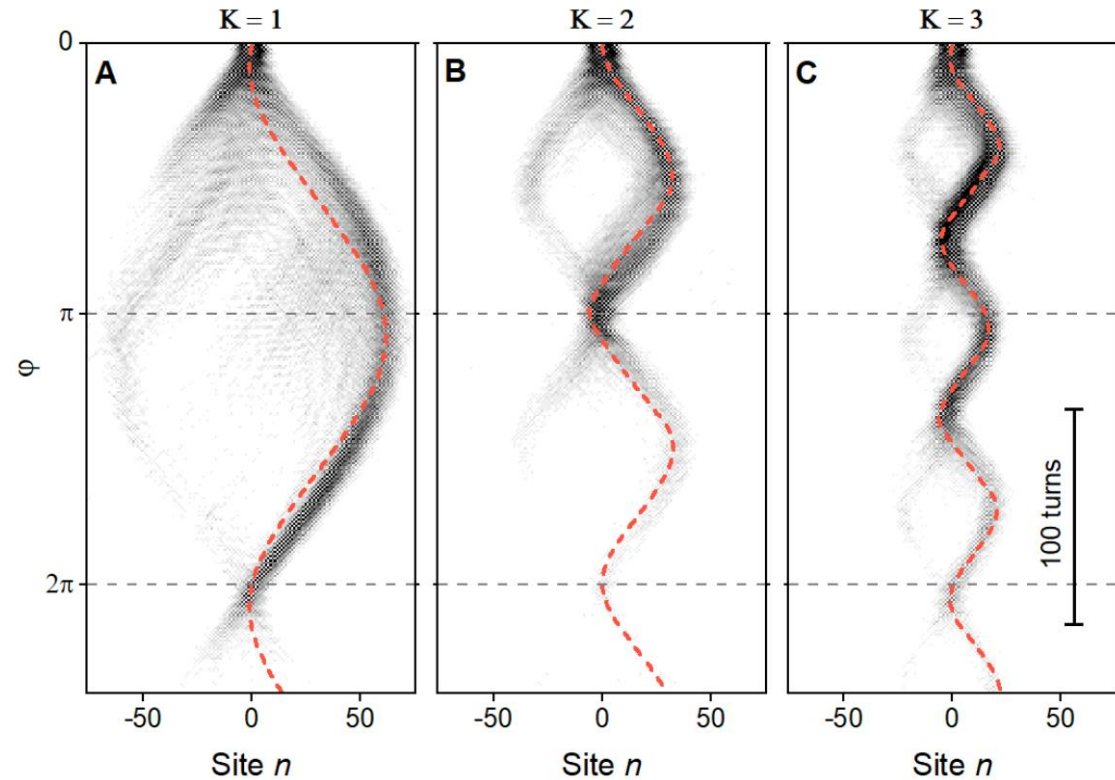
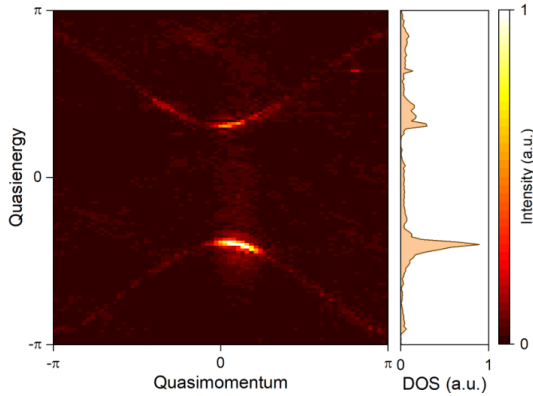
$K$  is a topological invariant:

$$\nu = \frac{1}{2\pi i} \int_0^{2\pi} d\varphi \operatorname{Tr} \left[ U_F^{-1} \frac{\partial U_F}{\partial \varphi} \right] = \sum_{j=\pm} \frac{1}{2\pi} \int_0^{2\pi} d\varphi \frac{\partial E_j}{\partial \varphi} = 2K$$





Exciting a single band:



Group velocity:

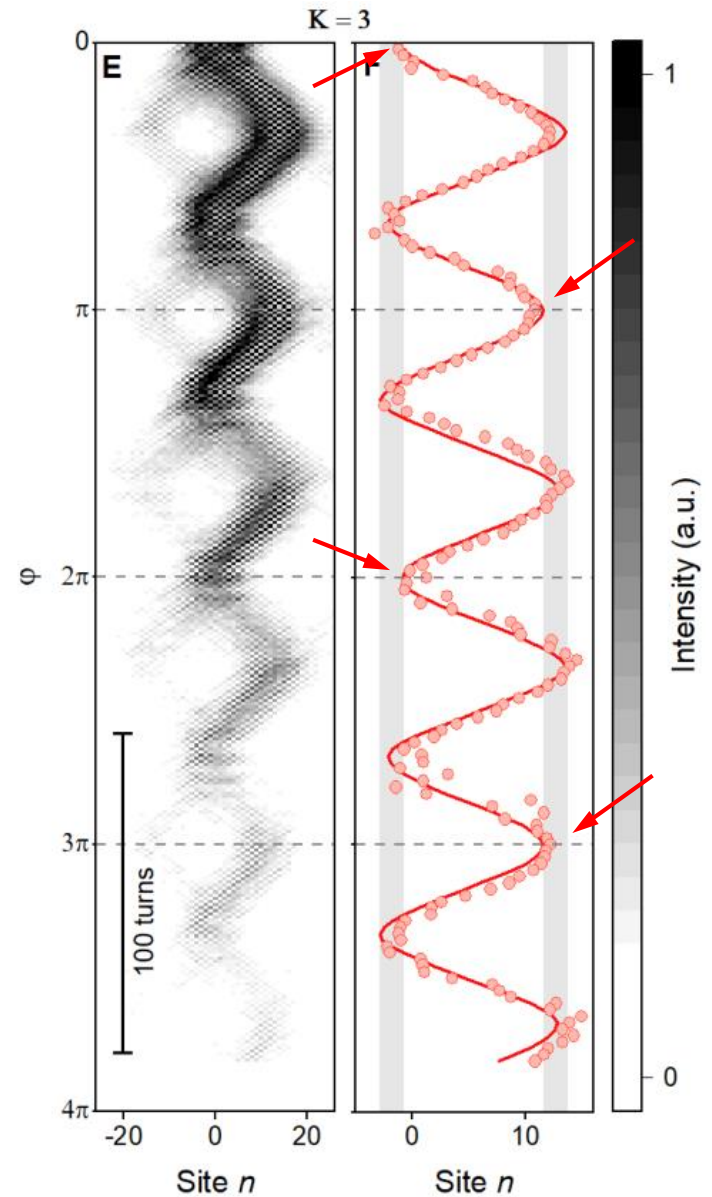
$$v_g^\pm(k, \varphi) = \frac{\partial E_\pm(k, \varphi)}{\partial k} = \pm \frac{\cos \theta_1 \cos \theta_2 \sin(k + K\varphi)}{\sqrt{1 - [\cos \theta_1 \cos \theta_2 \cos(k + K\varphi) - \sin \theta_1 \sin \theta_2 \cos(\Delta\varphi)]^2}}$$

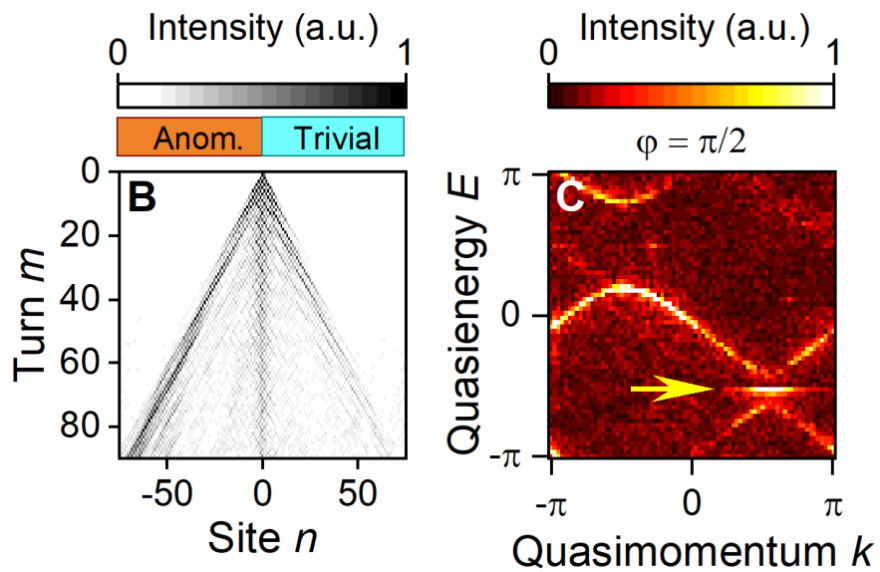
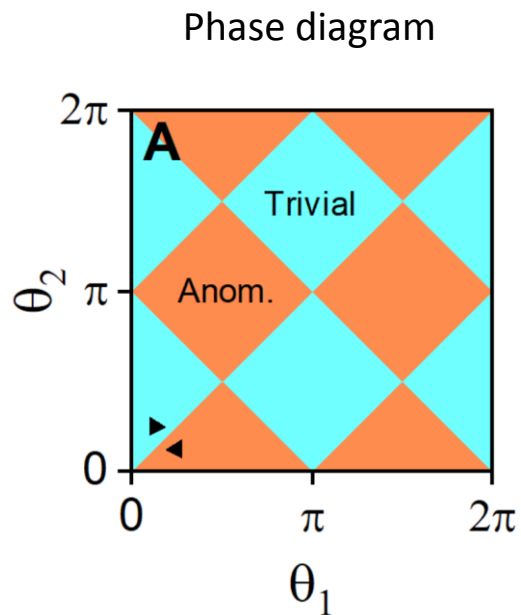
Frequency of sub-oscillations is topologically protected

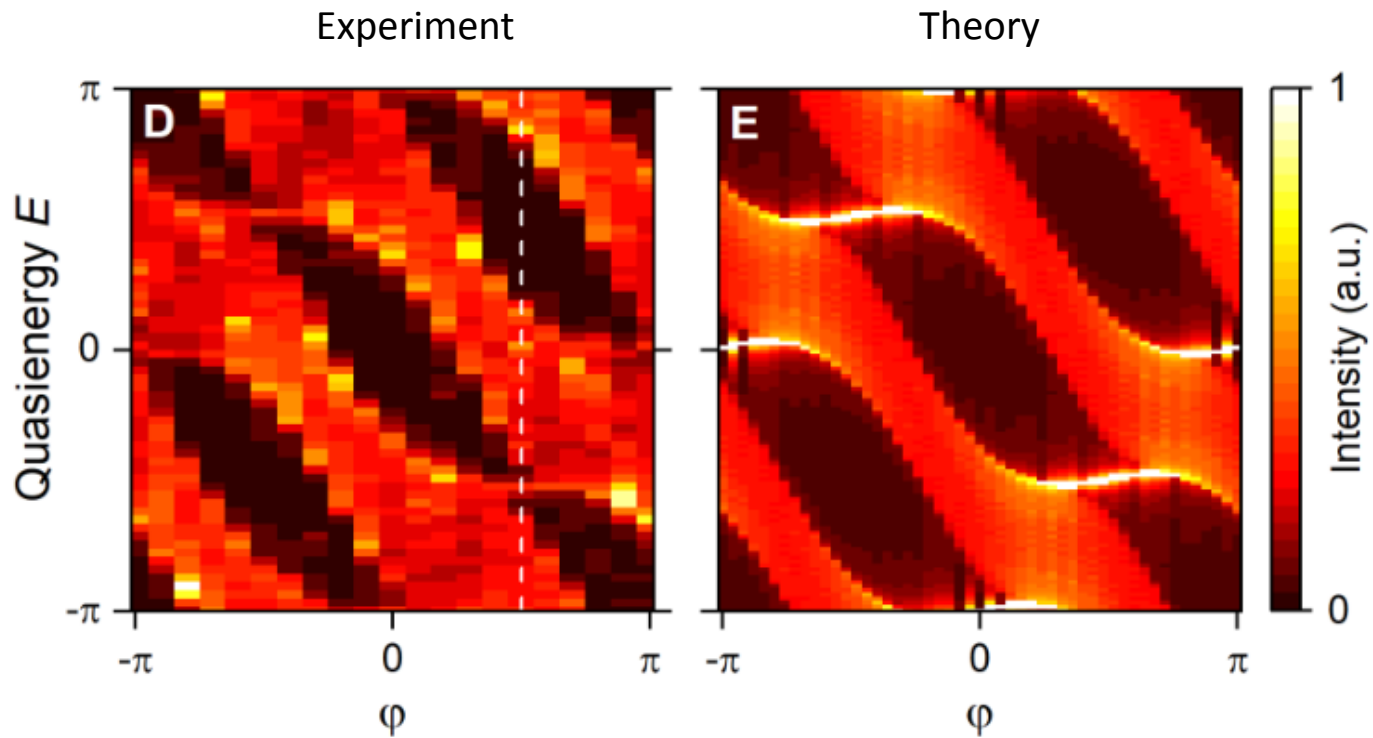
Amplitude of sub-oscillations can change!

Group velocity:

$$\pm \frac{\cos \theta_1 \cos \theta_2 \sin(k + K\varphi)}{\sqrt{1 - [\cos \theta_1 \cos \theta_2 \cos(k + K\varphi) - \sin \theta_1 \sin \theta_2 \cos(\Delta\varphi)]^2}}$$

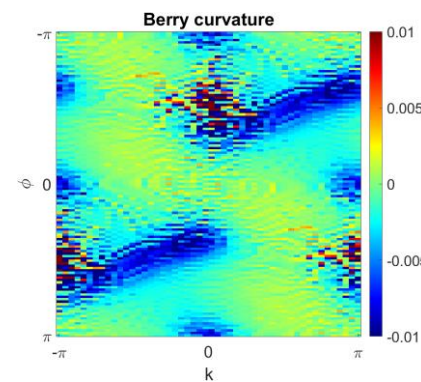
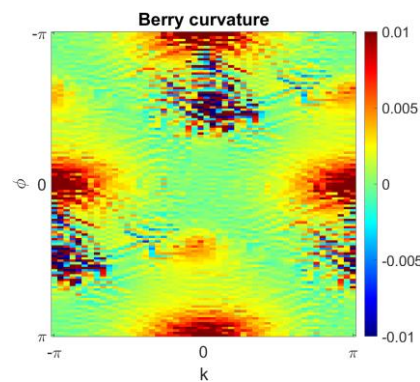
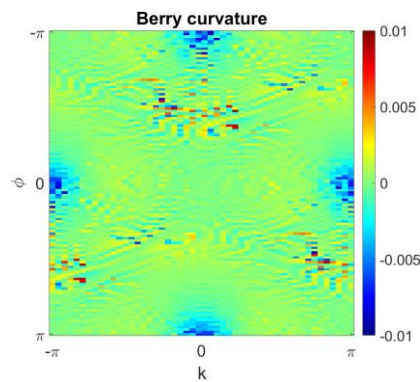
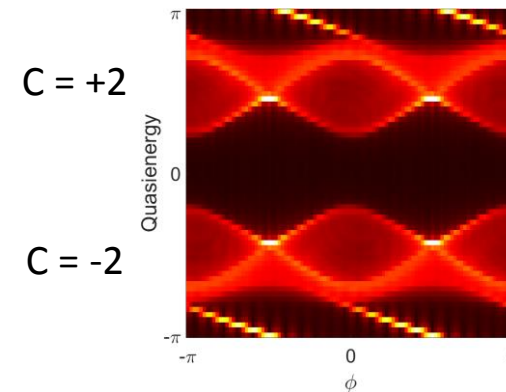
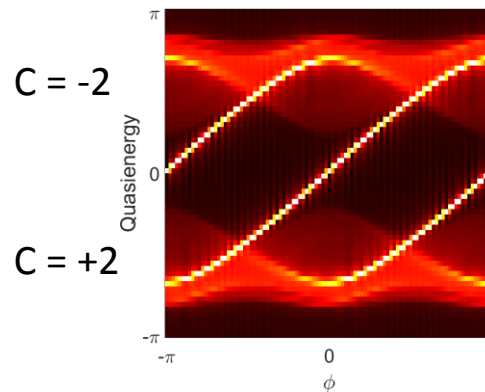
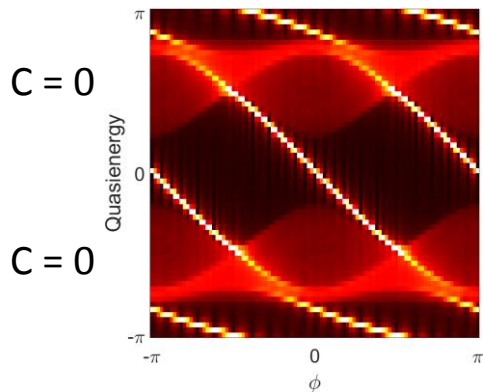
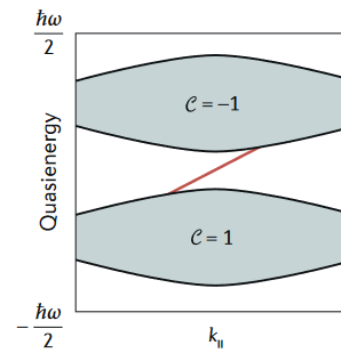
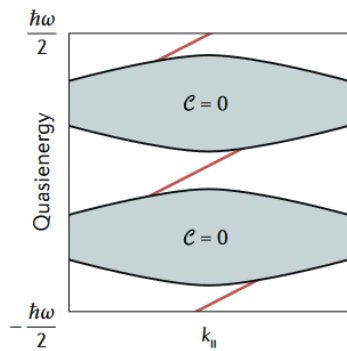




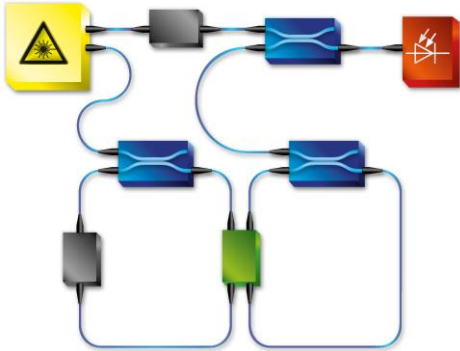


Two topological invariants coexist!

# Next steps: Chern number, Berry curvature

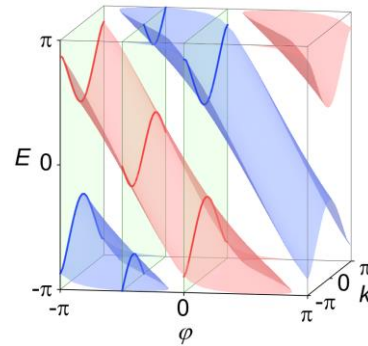


## Universal photonic simulator for studying topological effects



Lechevalier et al., Commun. Phys. 2021

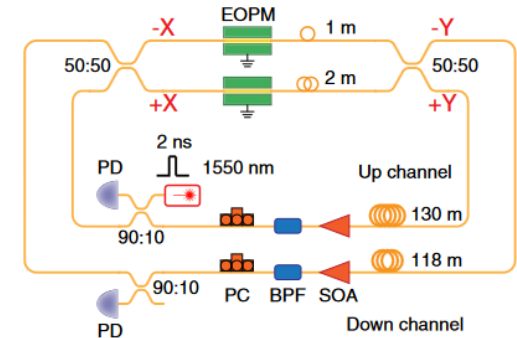
## Realization of a Floquet winding metal



Upreti et al., Phys. Rev. Lett. 2020  
Adiyattullin et al., arXiv:2203.01056

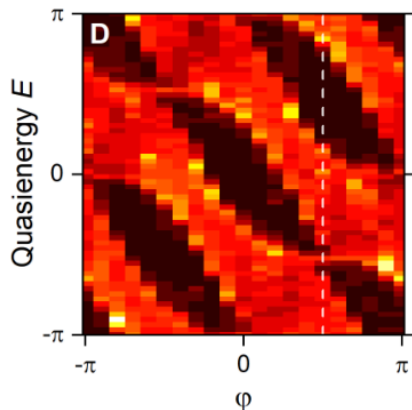
## Next steps

### Expanding to higher dimensions

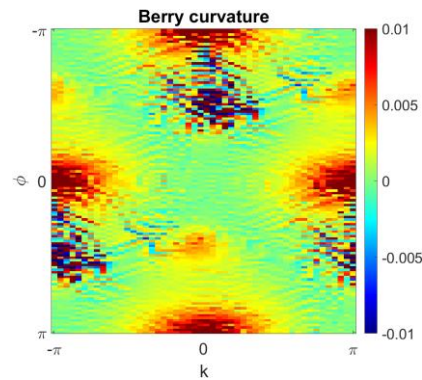


Chalabi et al., Phys. Rev. Lett. 123, 150503 (2019)

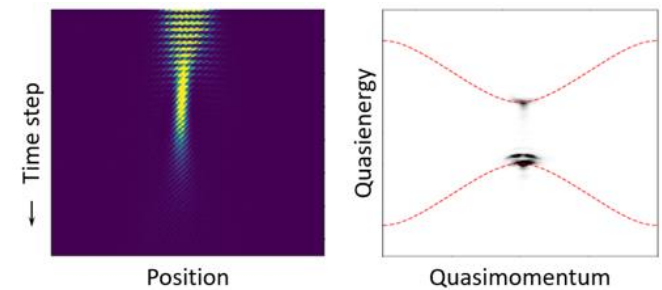
## Coexistence of two topological invariants



## Measuring the invariants

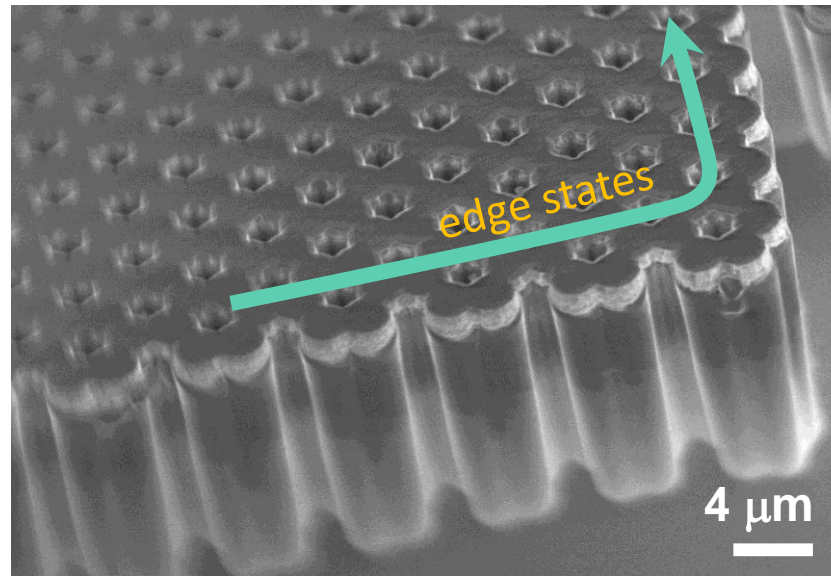


## Interactions due to fiber nonlinearities



by C. Lechevalier

2D lattice of polariton resonators



O. Jamadi et al., Light Sci. Appl. 9, 144 (2020)